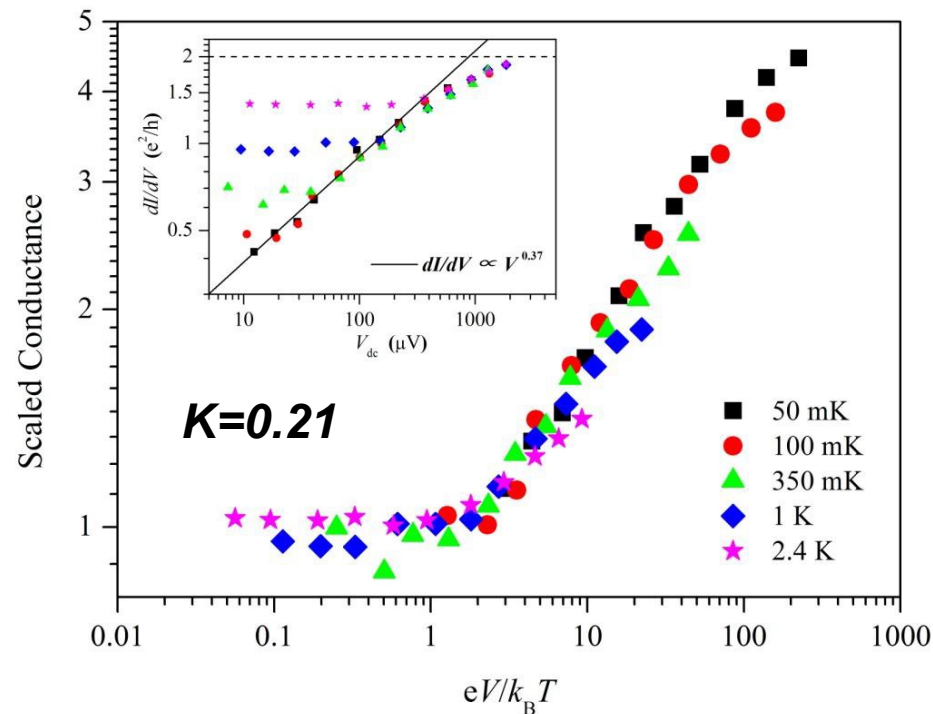
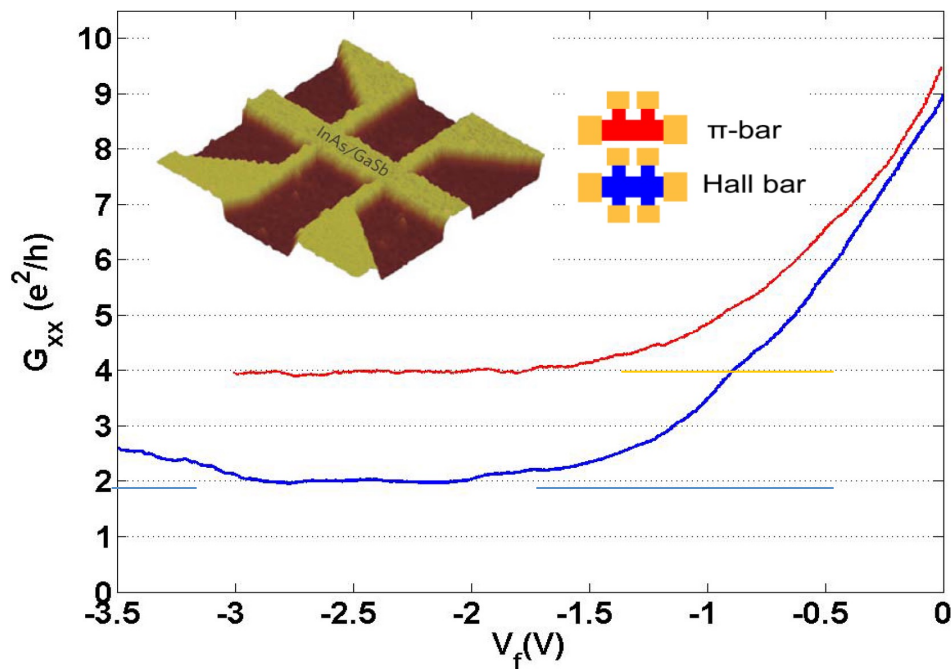


Gate-Controlled Helical Luttinger Liquid In InAs/GaSb Edges

Rui-Rui Du

Rice University/Peking University



ACKNOWLEDGE

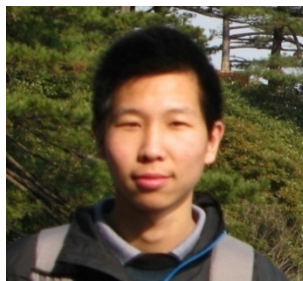
Rice

Ivan Knez

Lingjie Du

Peking U

Tingxin Li



Collaborations

Teledyne

Gerry Sullivan

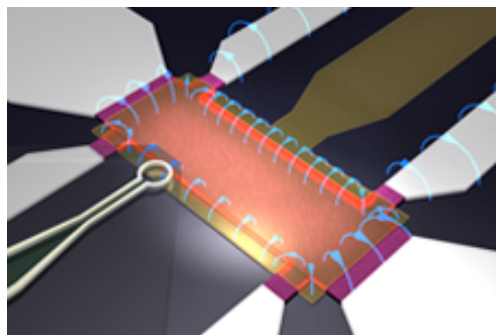
Stanford

Kathryn Moler

Eric Stanton



MBE



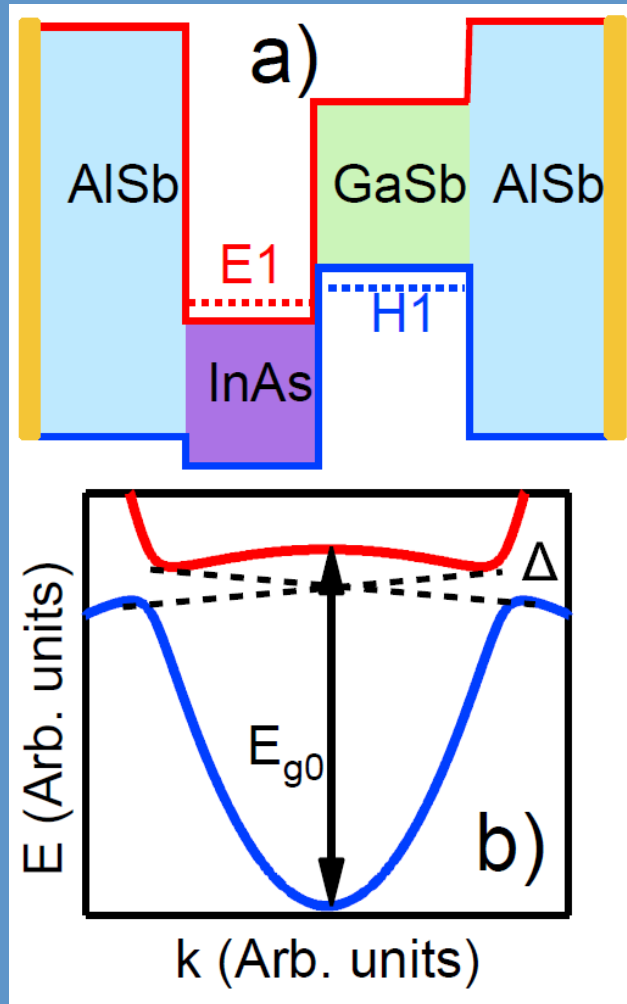
SQUID Imaging



OUTLINE

- 1. InAs/GaSb: Materials and Devices**
- 2. Quantized and Non-quantized Transport**
- 3. 1D- Helical Luttinger Liquid Transport**
- 4. Outlook: Tuning the Interactions**

6.1 A Family



Liu *et al*, *PRL* 100, 236601 (2008)

1. Band parameter

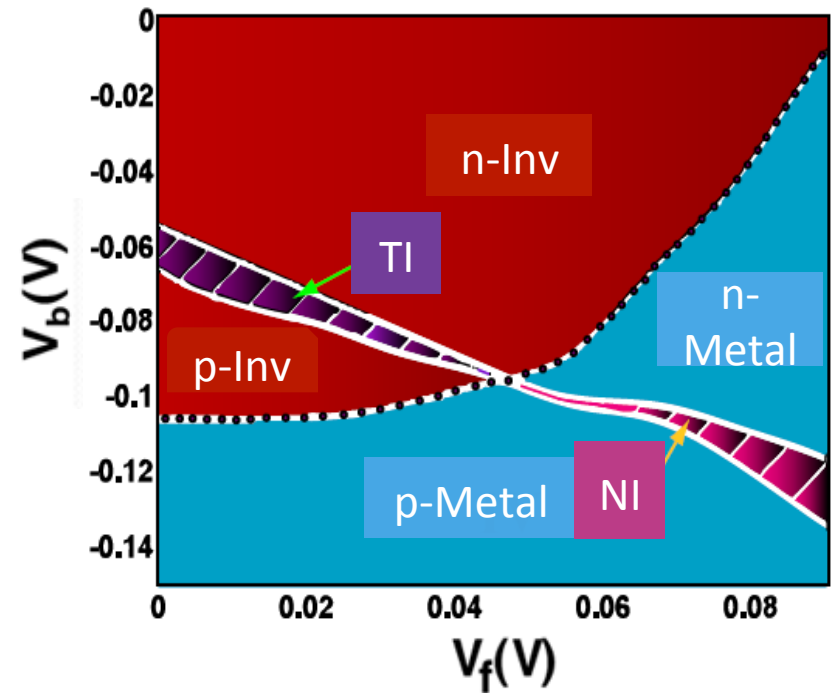
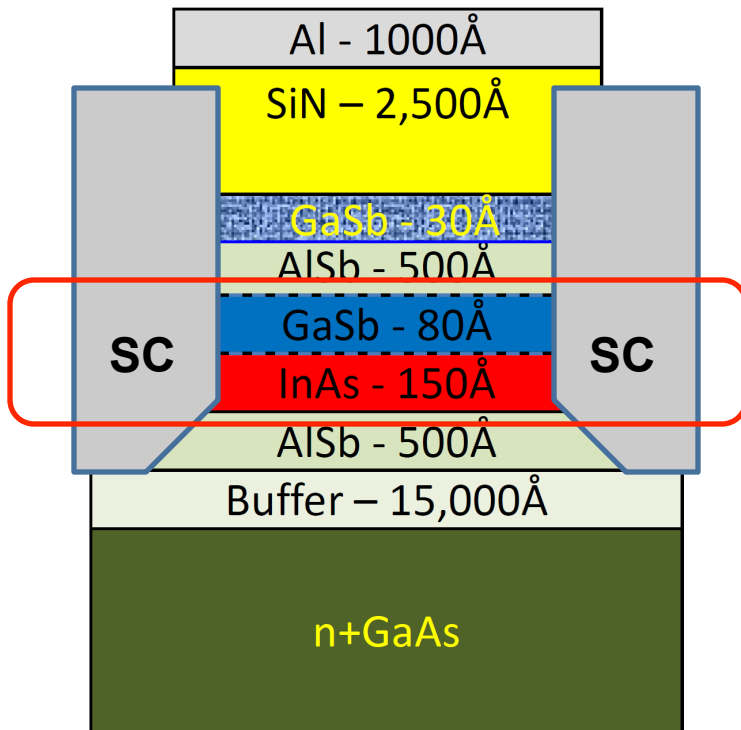
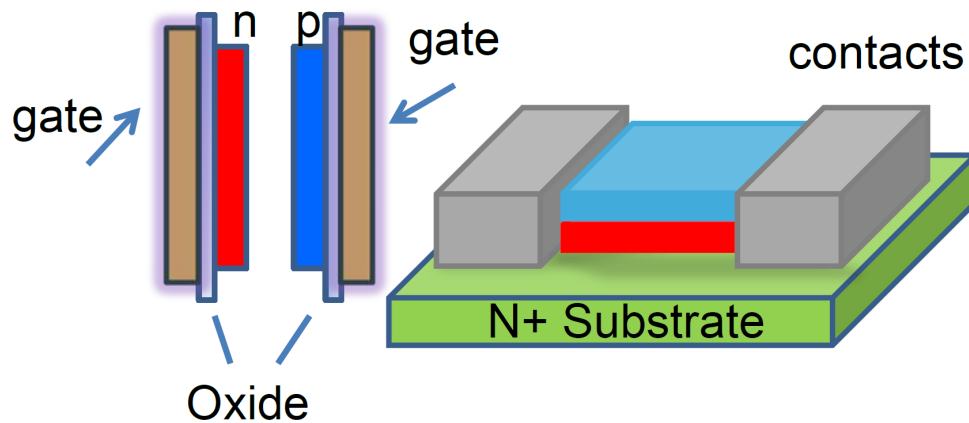
$$\Delta E = E1 - H1$$

tuned by width, x, gates

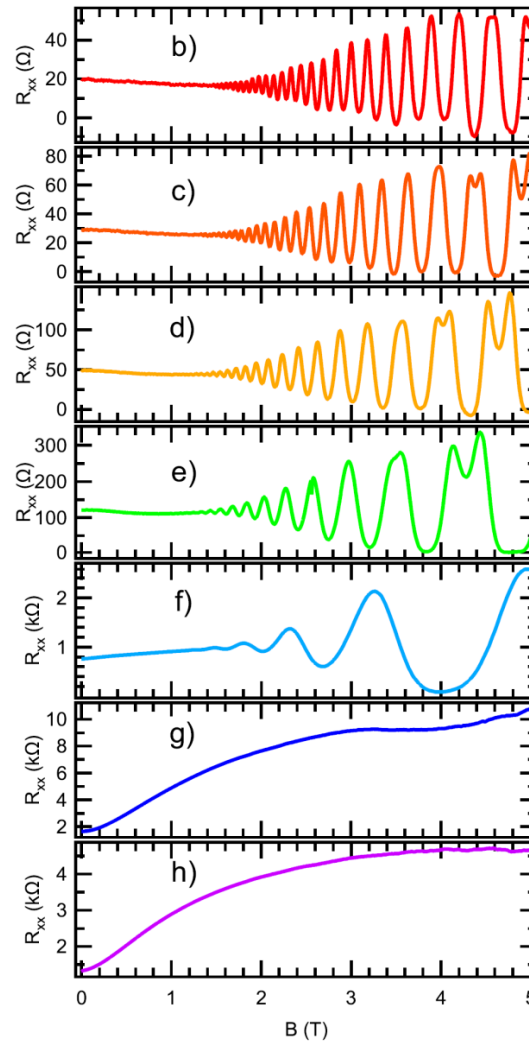
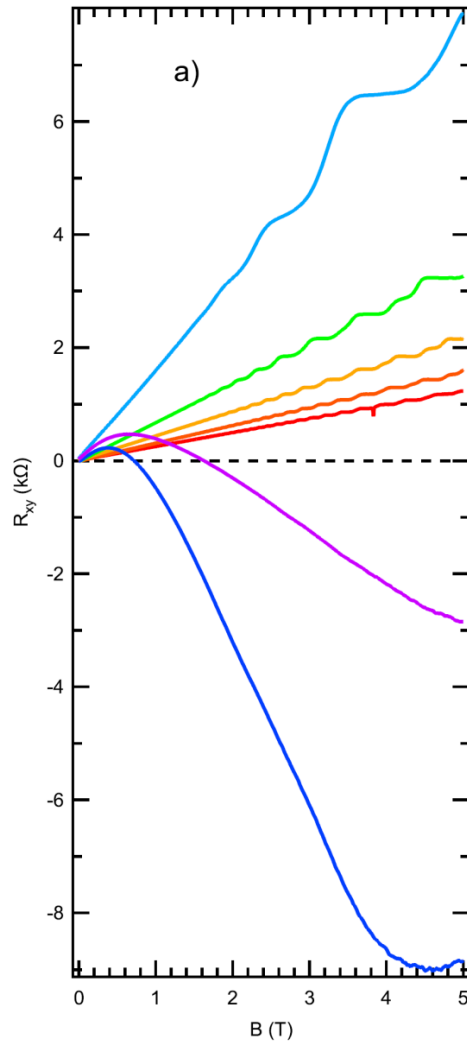
2. $\Delta E < 0$ inverted -0.15 eV to 0

Tunneling mix e-h and opens a gap at finite k -vector

Wafer and Device



General Quality



V front
= +6 V

+4

+2

0

-2

-4

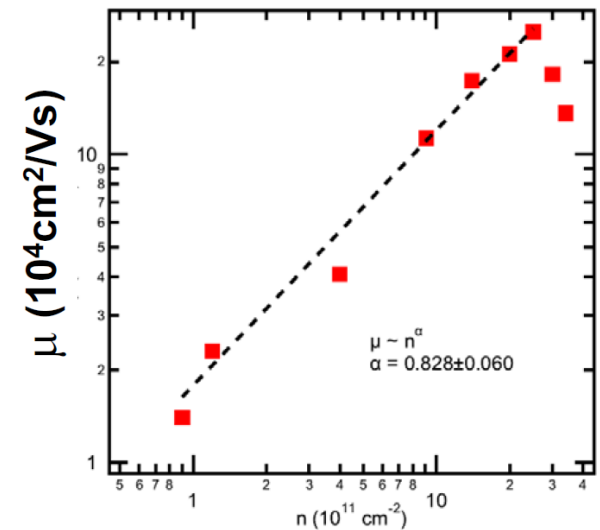
-6

Electron Mobility

$\sim 2 \times 10^5 \text{ cm}^2/\text{Vs}$

Hole mobility

$\sim 2 \times 10^4 \text{ cm}^2/\text{Vs}$

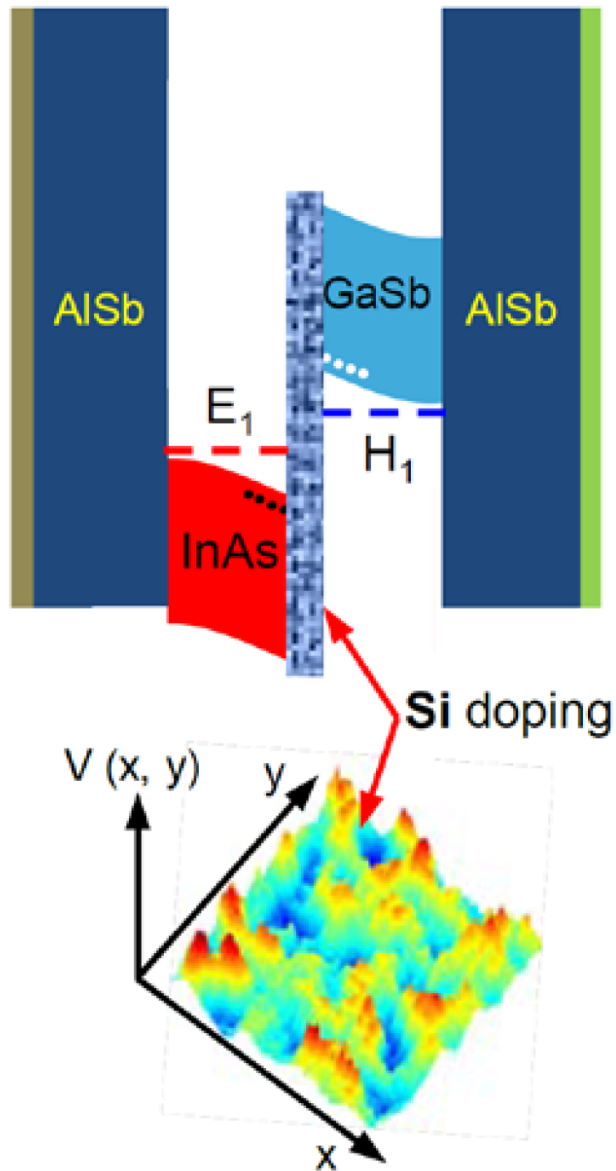


Ivan Knez, Ph. D thesis, Rice University 2011

<http://www.ruf.rice.edu/~dulab/index.html>

Quantized and Non-quantized Edge Transport

Si-doping at the interface

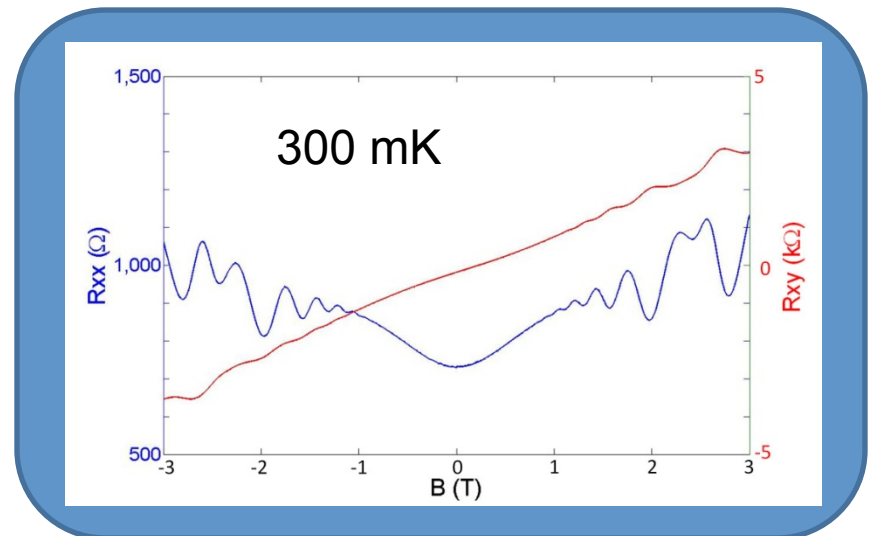


Si: donor in InAs, acceptor in GaSb

*Localize residual e and h

*Less effective for hybridized e - h

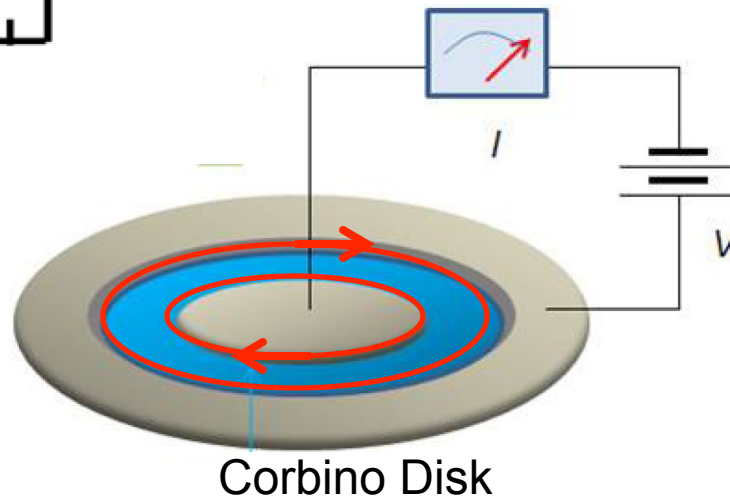
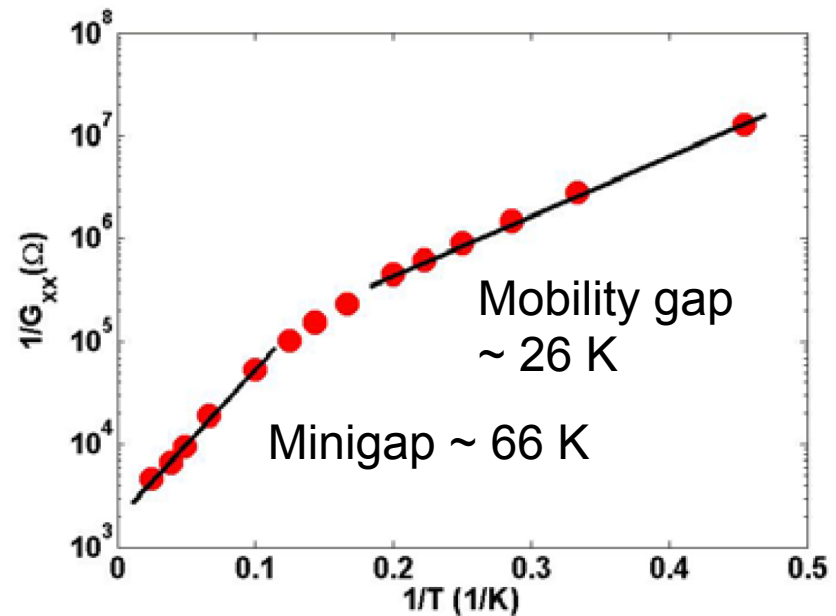
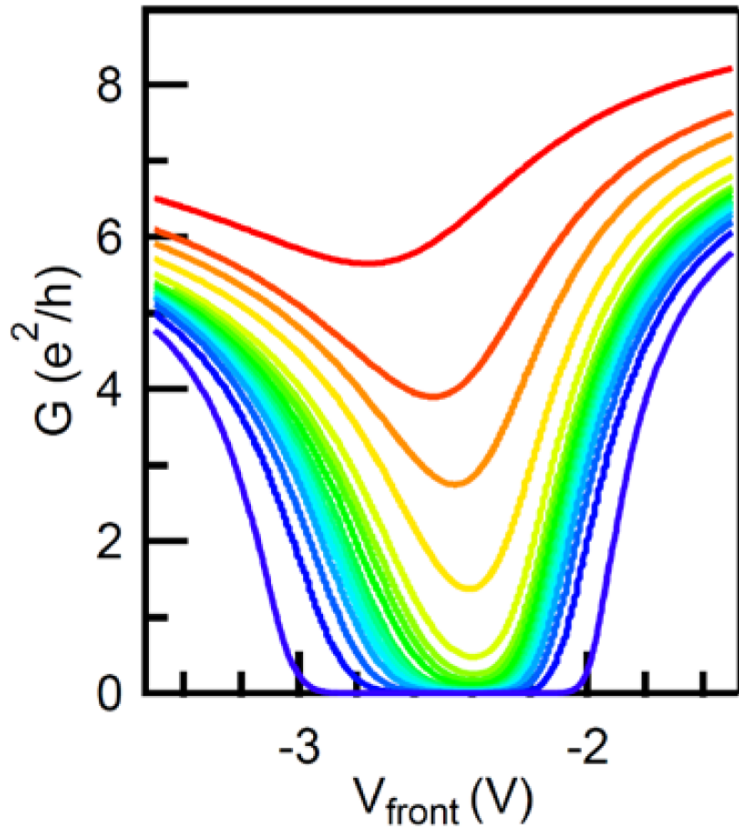
* Dipole field opposite to carrier field



40,000 cm^2/Vs

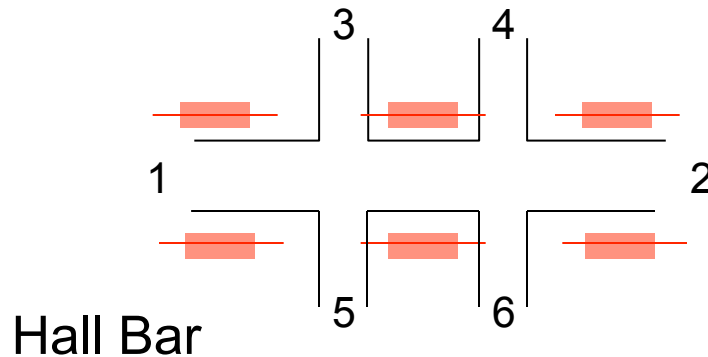
Truly Insulating Bulk

$\rho_{\text{Bulk}} > 100 \text{ M}\Omega/\text{Sqr}$
4 order of mag.

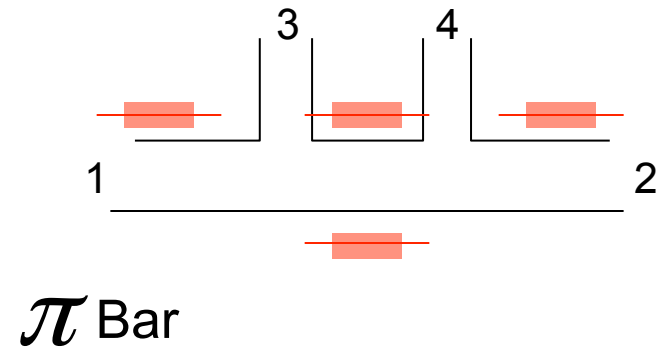


Conductance of Helical Edge Modes: Nonlocal - Transport

- For phase coherent sample Landauer-Buttiker formula gives conductance value



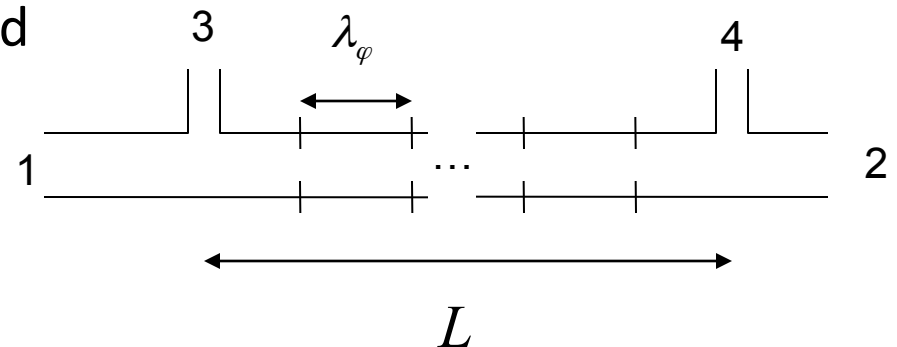
$$G_{12,34} = \frac{2e^2}{h}$$



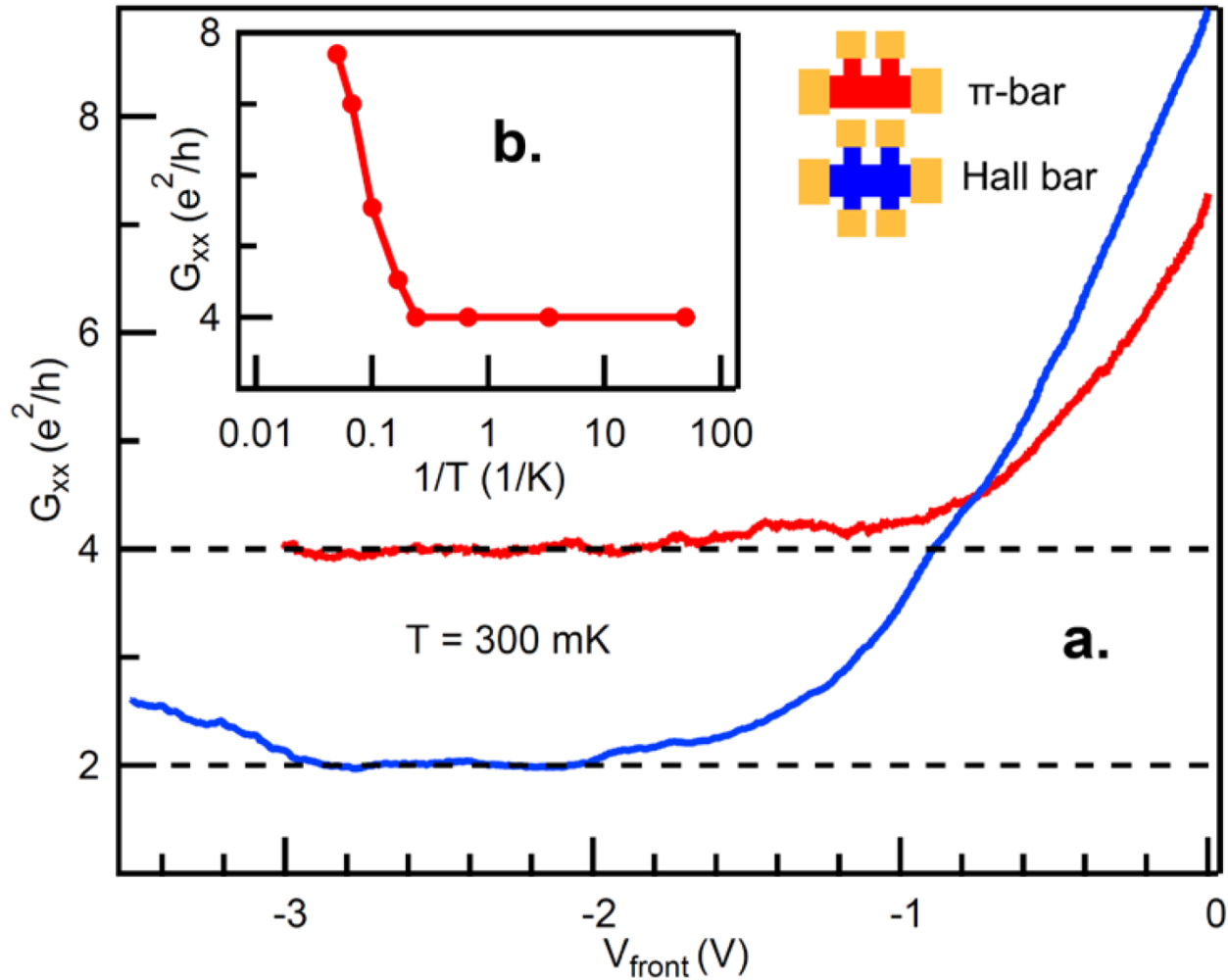
$$G_{12,34} = \frac{4e^2}{h}$$

- Longer samples can be modeled by inserting phase breaking probes and applying Buttiker formula:

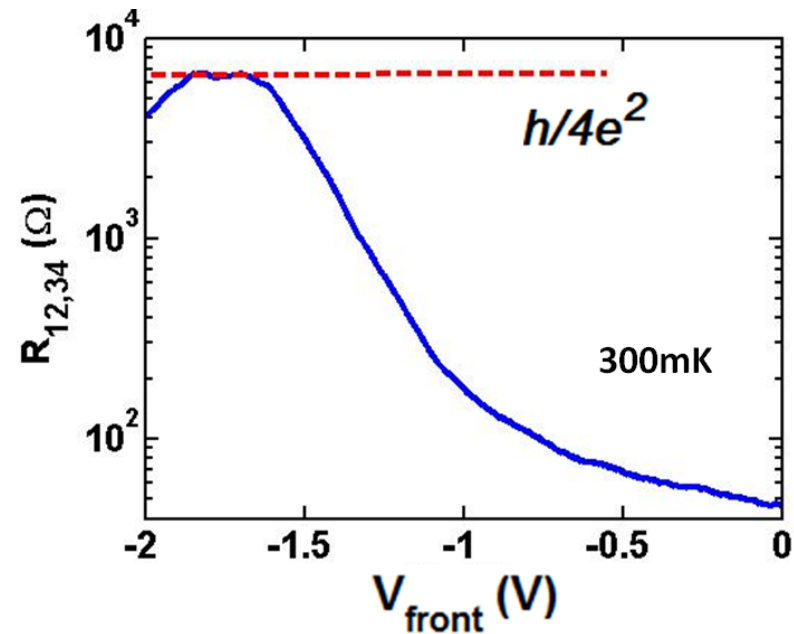
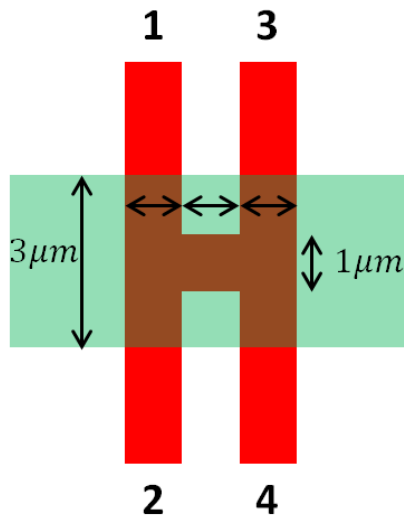
$$G_{12,34} = \frac{2e^2}{h} \left(\frac{\lambda_\varphi}{L} + \left(\frac{\lambda_\varphi}{L} \right)^2 \right)$$



Conductance Plateau up to 4 K

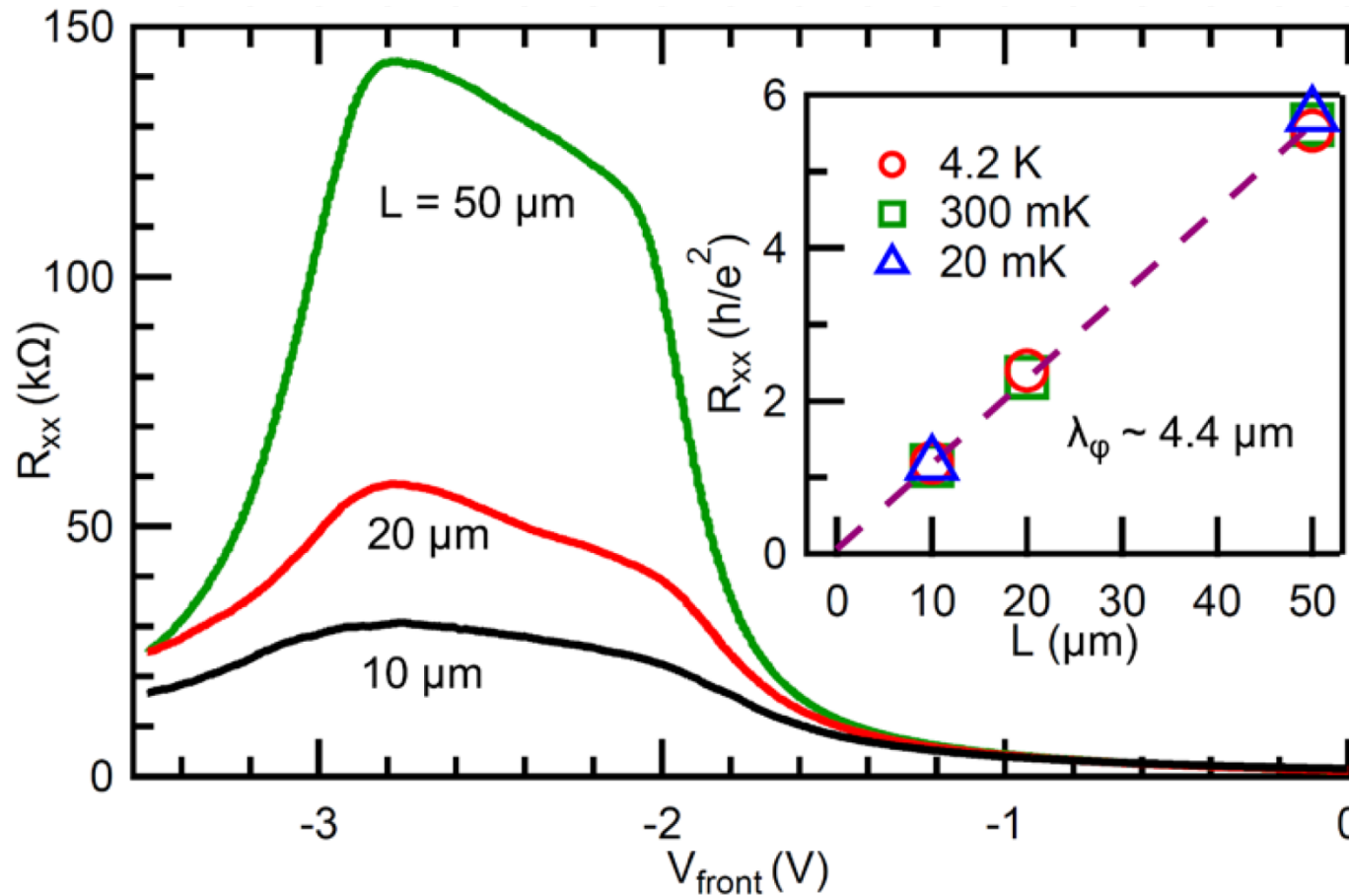


Non-Local Edge Transport: H-Bar



Lingjie Du

Longer Samples



$$G = \frac{2e^2}{h} \left(\frac{\lambda_\phi}{L} + \left(\frac{\lambda_\phi}{L} \right)^2 \right)$$

“Mean-Free-Path” $\lambda_\phi \sim 4 \mu m$
 Indep. with T 20 mK 4K

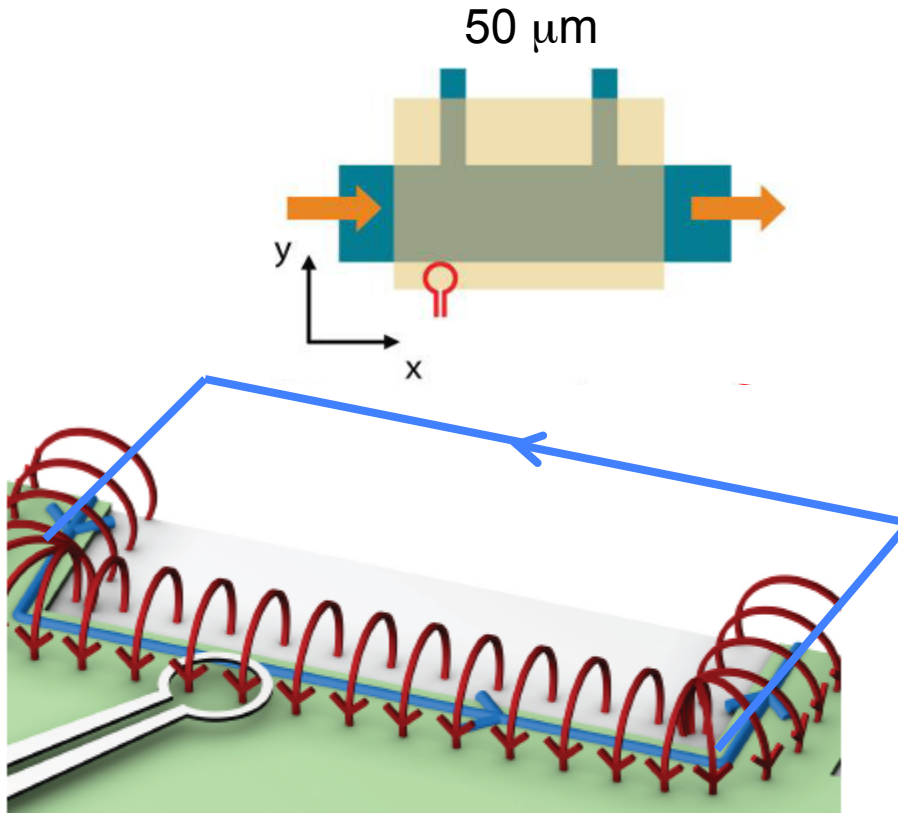


SQUID Imaging of InAs/GaSb Edge Modes

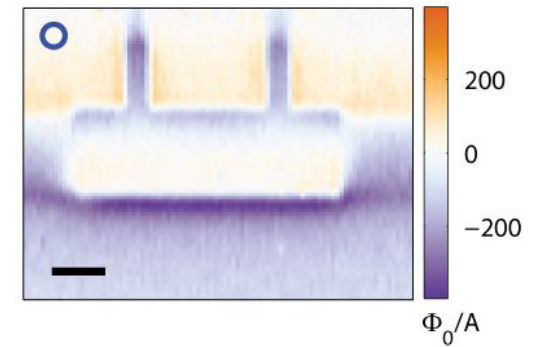
Stanford: Spanton, Moler

Rice: Lingjie Du, RRD

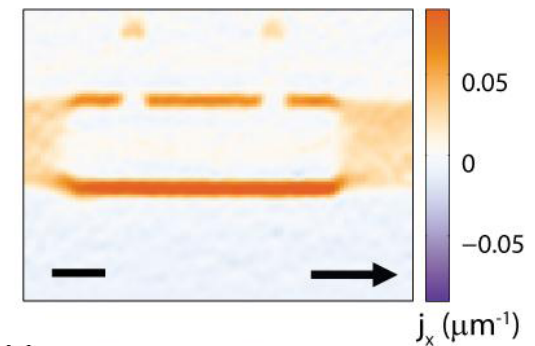
PRL 14



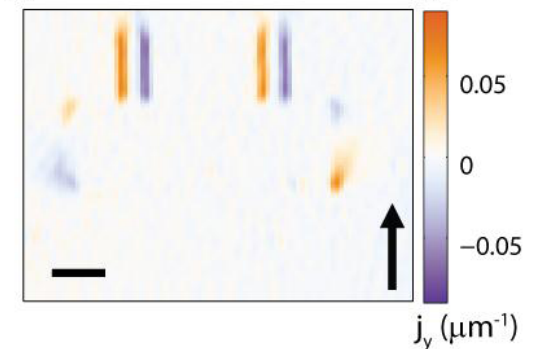
Flux



Current
(Horizontal)

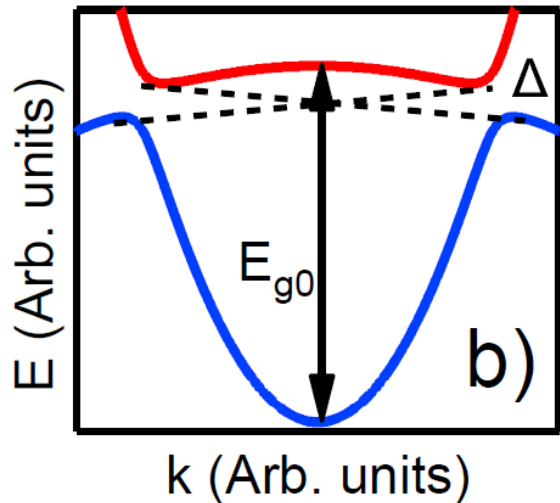


Current
(Vertical)



Properties of InAs/GaSb Edge States

- Quantized plateau for short samples (a few μm)
- Non-quantized conductance for long samples
- Conductance nearly T-independent over large range of T
- Robust under high B



- Low velocity of edge modes

$$v_F \sim 2 \times 10^4 \text{ m/s}$$

$$k_F \sim 1 \times 10^6 / \text{cm}$$

- Extreme. long scattering time

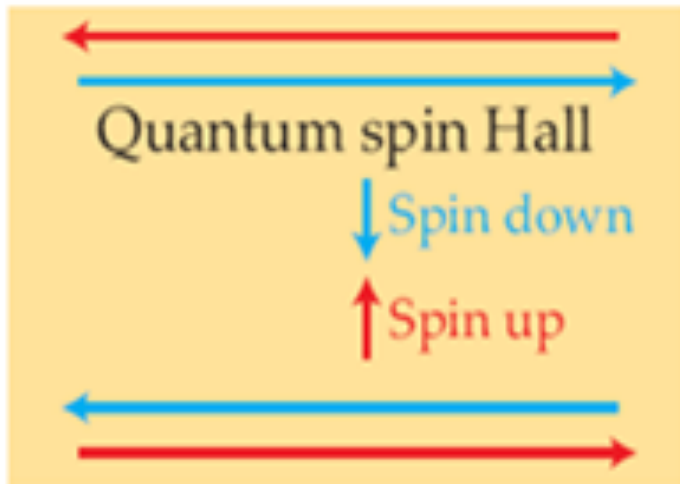
$$\tau = \lambda / v_F = 4 \mu\text{m} / (2 \times 10^4 \text{ m/s}) \sim 200 \text{ ps}$$

- Origins of back-scattering?

- Tunable mode dispersion

Interaction Effects in InAs/GaSb Helical Edge States

Non-Interacting Model:



1. TRS protected topological Insulator
2. Spin-momentum locking
3. non-magnet. impurity does not cause backscattering

Non-linear Edge Channel Conduct.

C. L. Kane and M. P. A. Fisher, Transport in a one-dimensional Luttinger liquid, *Phys. Rev. Lett.* **24**, 1220 (1992); *Phys. Rev. B* **46**, 15233 (1992)

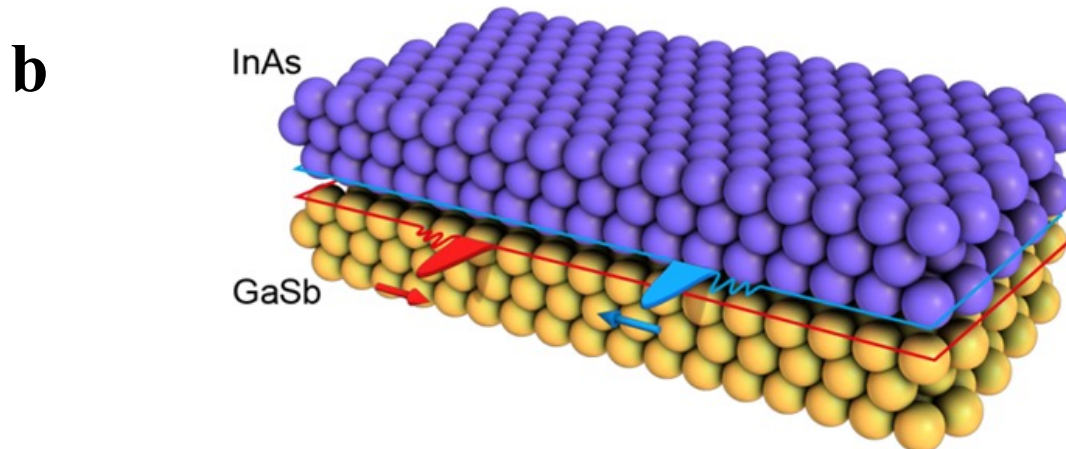
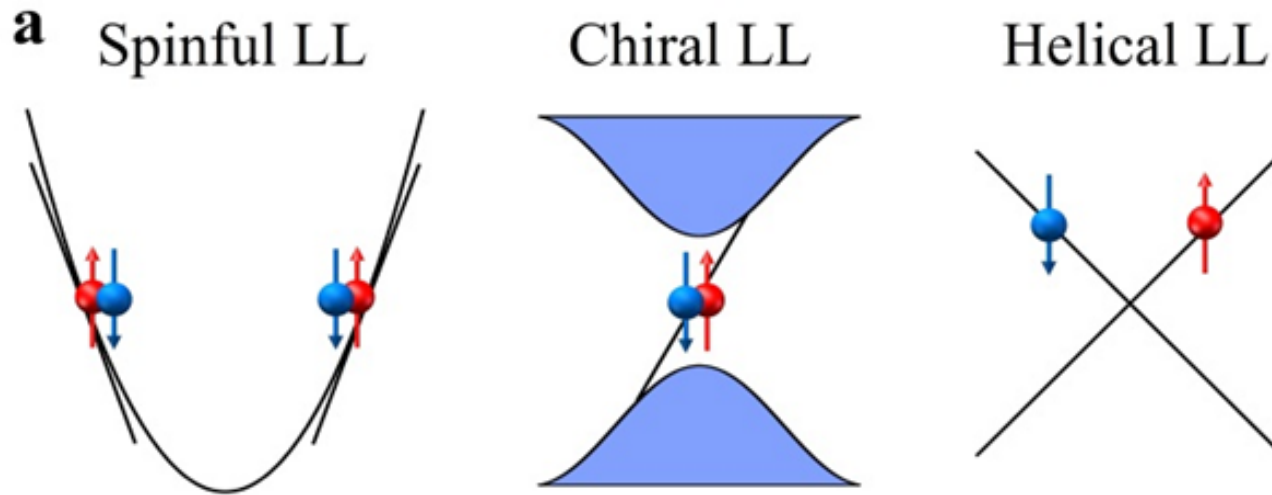
M. P. A. Fisher and L. Glazman, Transport in a one-dimensional Luttinger Liquid.

Wu, C., Bernevig B. A. & Zhang S. C. Helical liquid and the edge of quantum spin Hall systems. *Phys. Rev. Lett.* **96**, 106401 (2006).

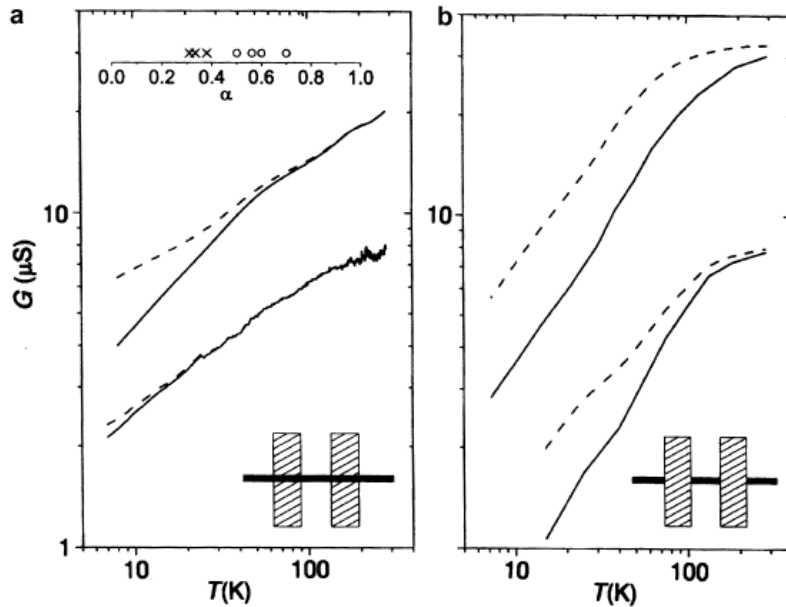
Teo, J. C. Y. & Kane, C. L. Critical behavior of a point contact in a quantum spin Hall insulator. *Phys. Rev. B* **79**, 235321 (2009).

Maciejko, J. *et al.* Kondo effect in the helical edge liquid of the quantum spin Hall state. *Phys. Rev. Lett.* **102**, 256803 (2009).

Family of Luttinger Liquids

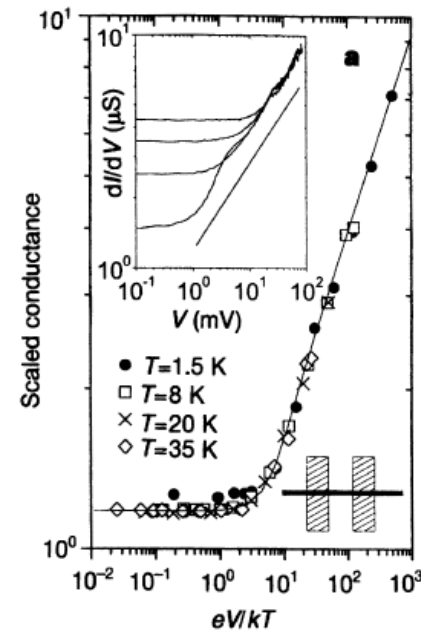


Carbon Nanotube (Spin-Full LL)

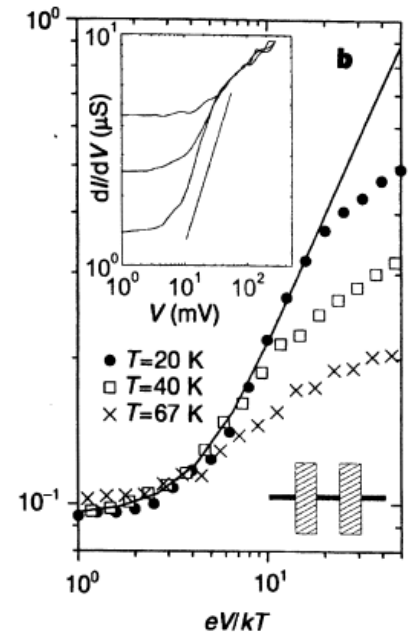


"Bulk Contacted"

"End Contacted"



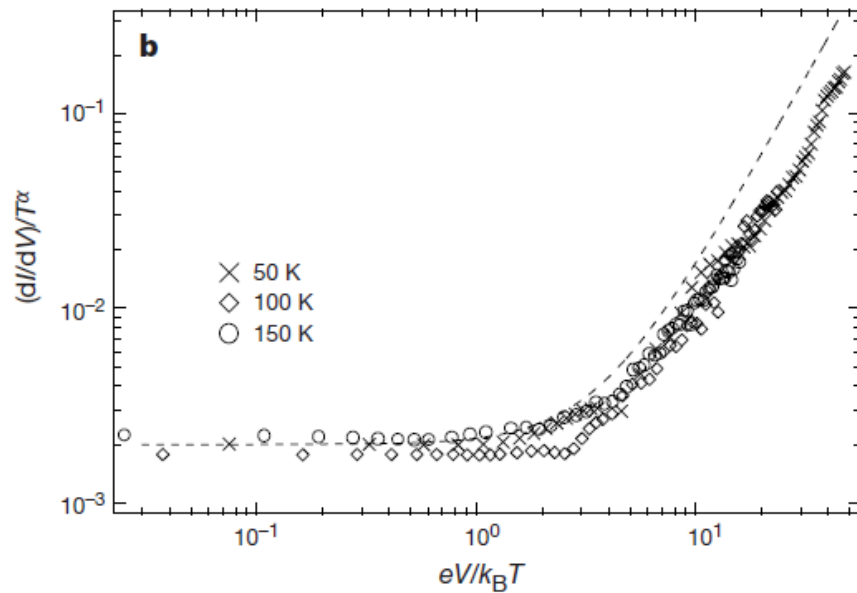
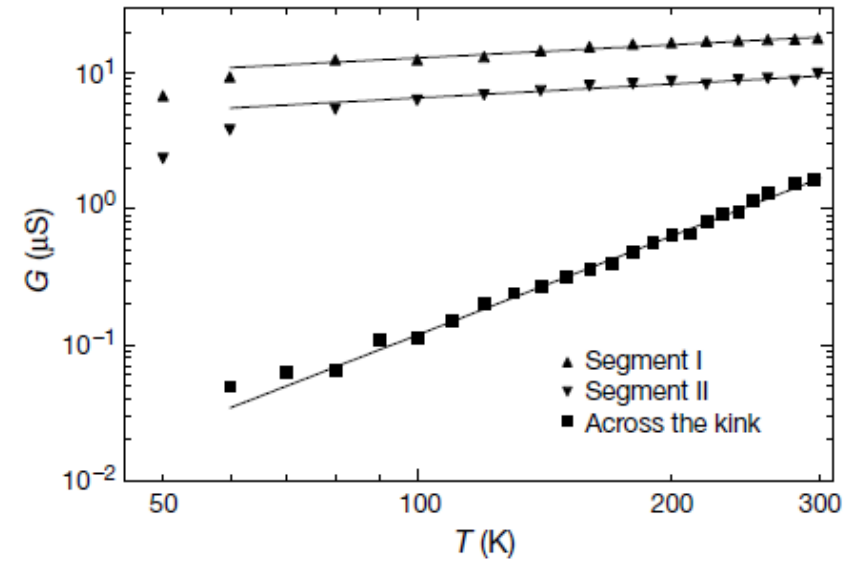
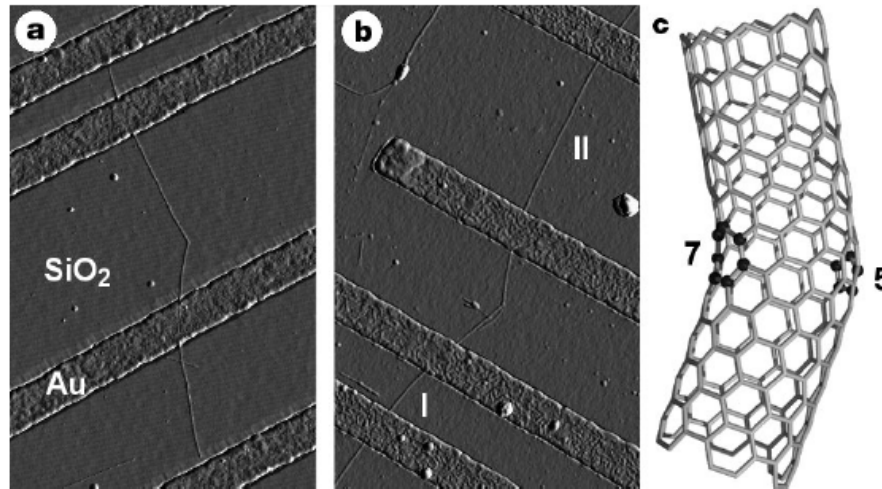
"Bulk Contacted"



"End Contacted"

- Metal contacts connect to rope of carbon nanotubes, serve as tunneling probes.
- $G \propto T^\alpha$, $dI/dV \propto V^\alpha$; $\alpha_{\text{bulk}} \sim 0.3$, $\alpha_{\text{end}} \sim 0.6$

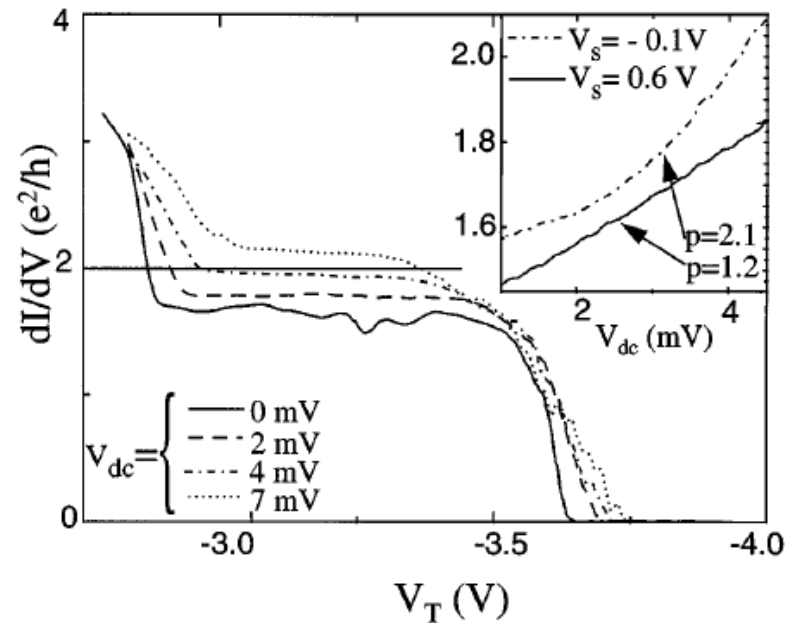
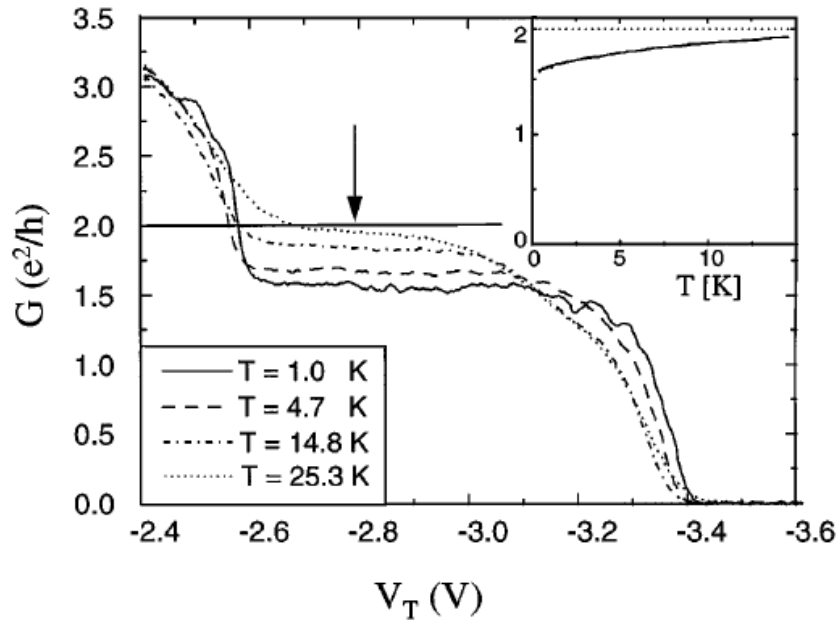
Carbon Nanotube (Spin-Full LL)



- Transport through Single-wall carbon nanotubes (SWNTs) with a kink.
- $G \propto T^\alpha$, $dI/dV \propto V^\alpha$

Z. Yao *et al.*, *Nature* **402**
273 (1999).

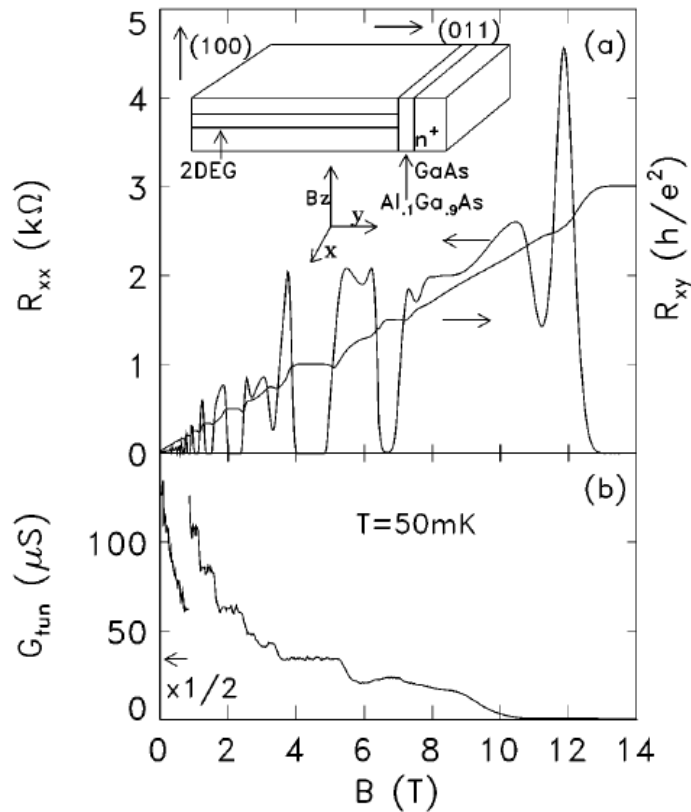
Quantum wires (Spin-Full LL)



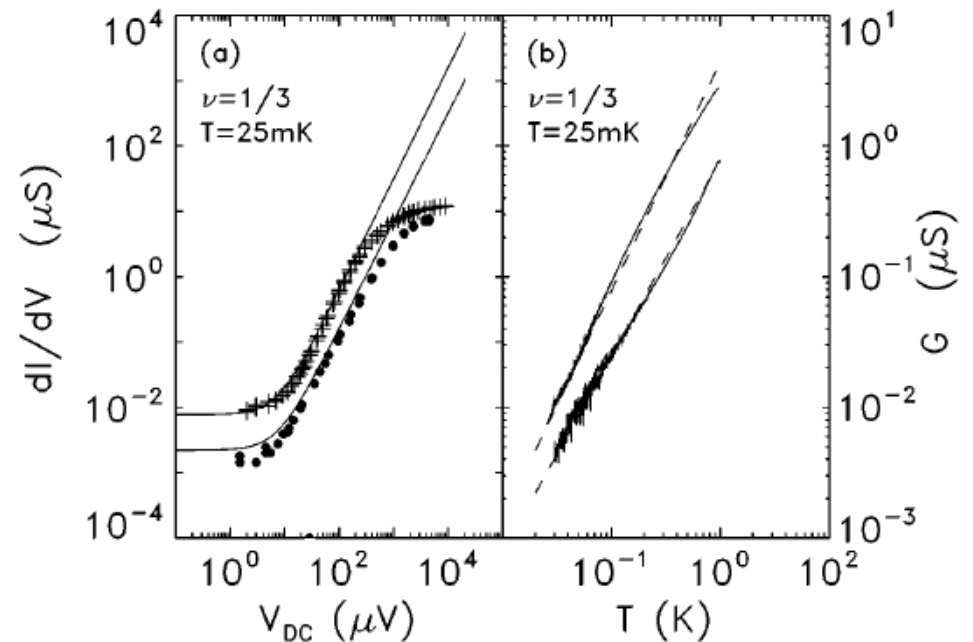
- High-quality quantum wires fabricated in GaAs-AlGaAs by using cleaved edge overgrowth.
- The conductance plateaus deviate from the universal quantized value at low temperature and low bias.
- $dI/dV = c + AV^p$

A. Yacoby *et al.*, *PRL* 77 4612
(1996).

Fractional Quantum Hall Edge (Chiral LL)



Cleaved-edge overgrowth



- $\nu = 1/3$ FQH edge
- $dI/dV \propto V^{1.7}$, $G \propto T^{1.75}$

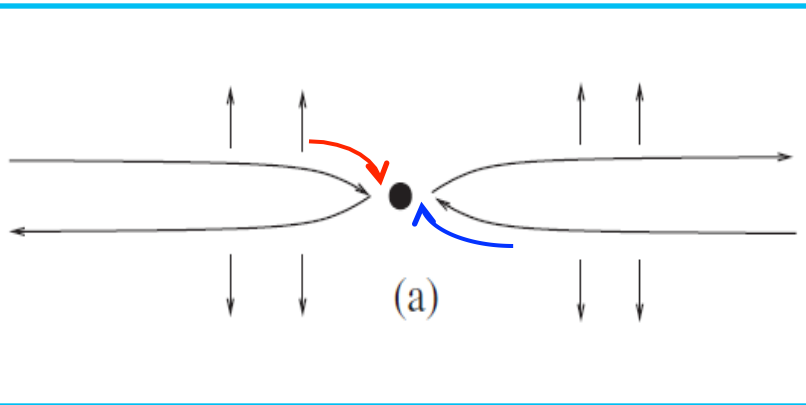
A. M. Chang *et al.*, *PRL* 77 2539
(1996).

Interacting Model: $1 > K > 1/2$; $1/2 > K > 1/4$;

Correlated 2-Particle Scattering is Allowed by TRS

C. Wu, B. A. Bernevig, and S. C. Zhang, Phys. Rev. Lett. **96**, 106401 (2006).

C. Xu and J. E. Moore, Phys. Rev. B **73**, 045322 (2006).



**2 particle correlated
Backscattering for $k < 1/4$**

---Breaking into two segments

PHYSICAL REVIEW B **73**, 045322 (2006)

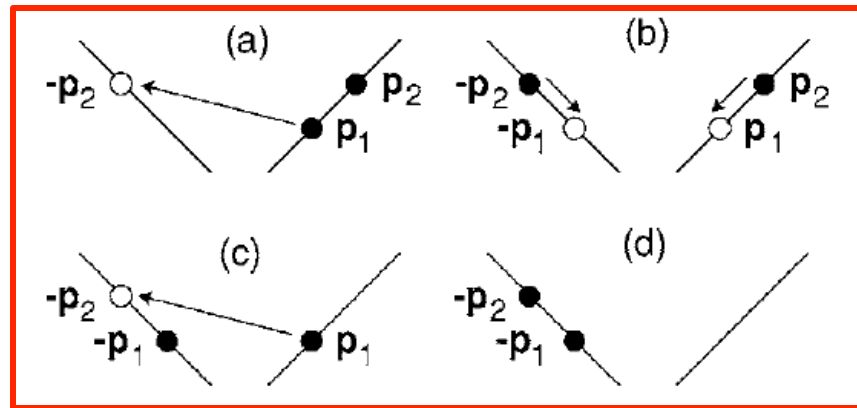


FIG. 1. An allowed two-particle backscattering process: (a) particle at momentum p_1 scatters to intermediate state $-p_2$; (b) particles at $\pm p_2$ interact and become intermediate states $\pm p_1$; (c) and (d), particle at intermediate state p_1 backscatters to state $-p_2$.

Calculating Luttinger Parameter K

PRB **79** 235321 (2009), and *PRL* **102** 256803 (2009):

$$K = \left[1 + \frac{2}{\pi^2} \frac{e^2}{\varepsilon \hbar v_F} \ln \left(\frac{7.1d}{\xi + 0.8w} \right) \right]^{-1/2}$$

Where ε is the bulk dielectric constant,

v_F is Fermi velocity of edge state,

d is the distance from the QW layer to a nearby metallic gate as a screening length for the Coulomb potential,

$\xi = 2\hbar v_F / E_{gap}$ (E_{gap} is the gap of bulk QSHI)

w is the thickness of the QW layer.

***HgTe/CdTe* Quantum Wells**

$K = 0.8$

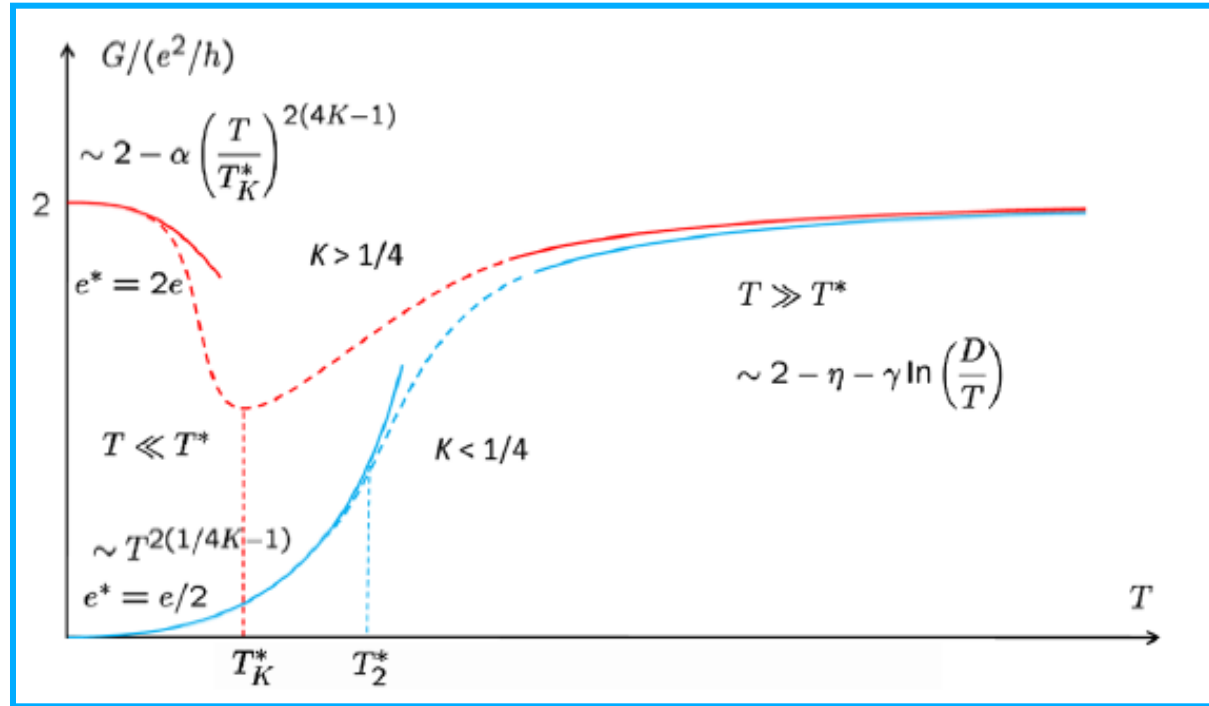
***Gated InAs/GaSb* Bilayers**

$K = 0.2 \sim 0.25$

***Strained InAs/InGaSb* QWs**

$0.5 > K > 0.25$

Possible Interacting Effects



Maciejko et al, PRL 09

Teo and Kane, PRB 09

Helical Liquid with strong
Coulomb interactions:

1. Log T -dependence at high T
2. Insulating phase at low T

Relevant to InAs/GaSb

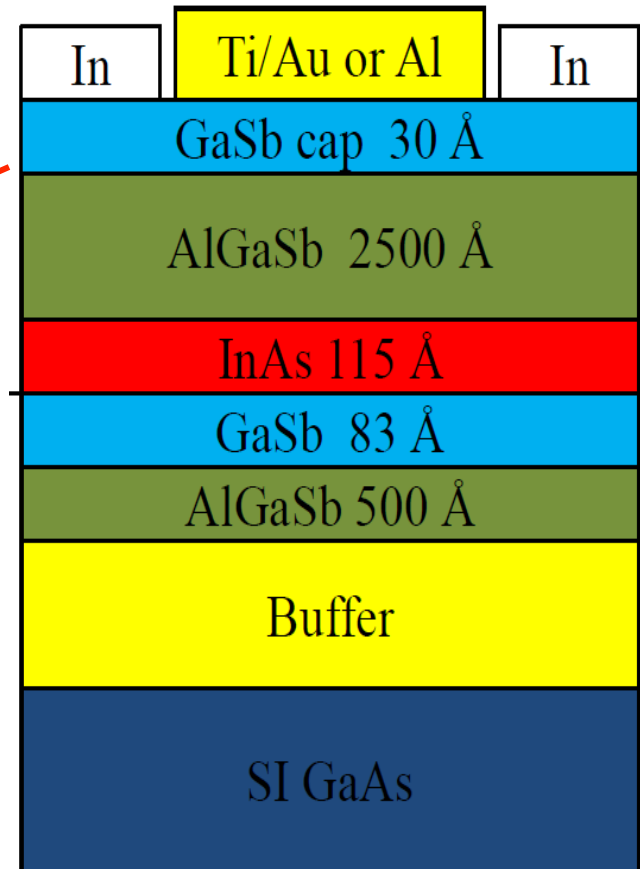
2 Samples for Ultralow T Measurements

Wafer A



$K \sim 0.23$

Wafer B

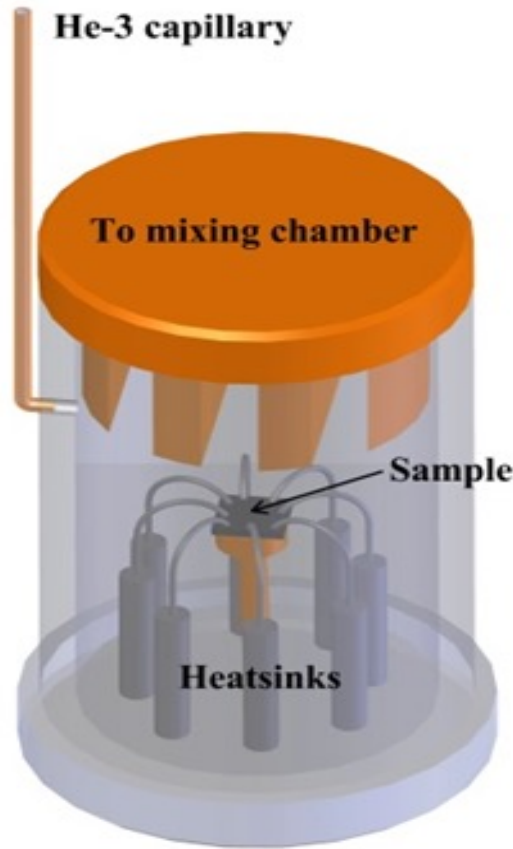


$K \sim 0.21$

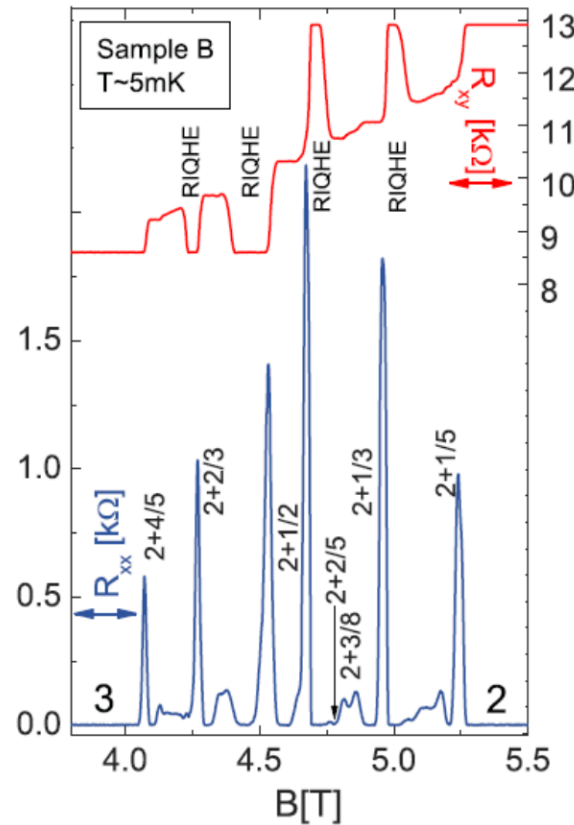
Cooling Electrons

Purdue Group: Gabor Csathy / Kate Schreiber

PKU group: Xi Lin/ Pengjie Wang, Hailong Fu



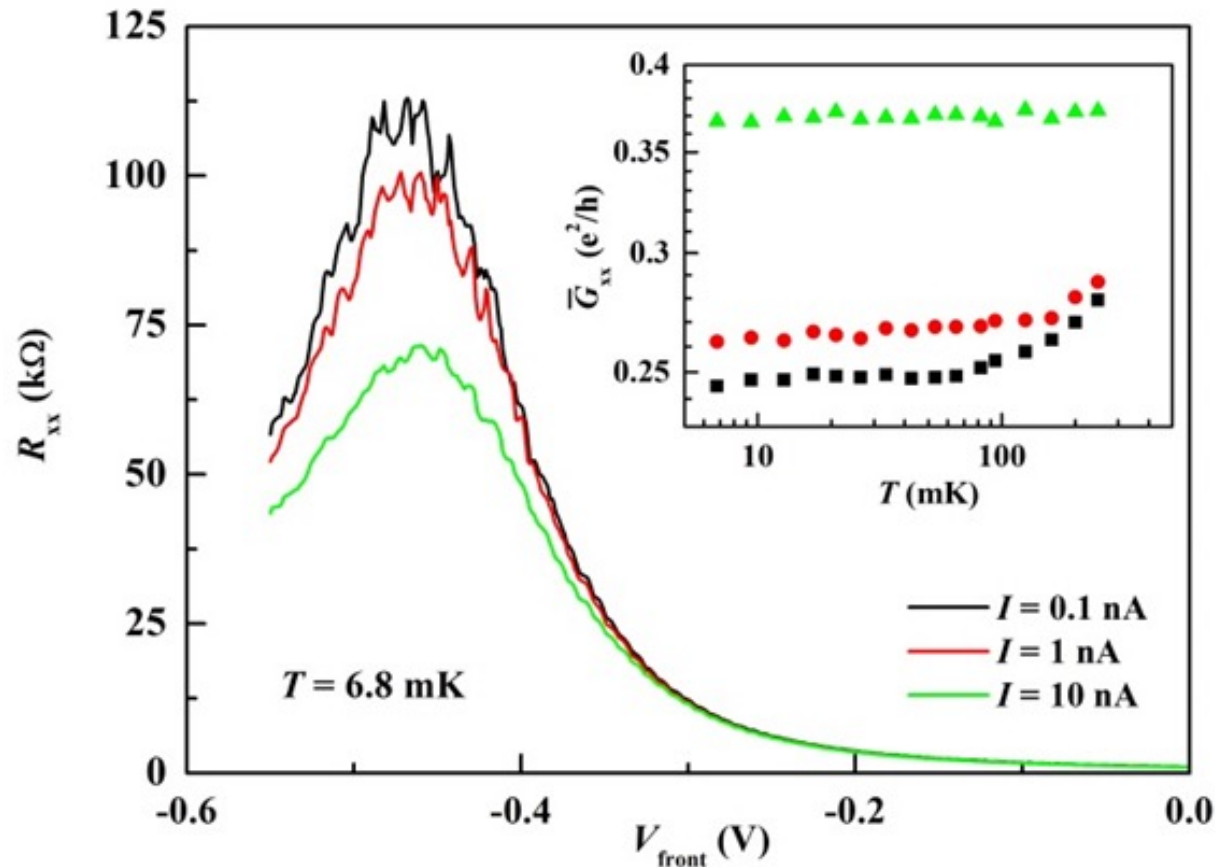
He3 Immersion Cell



5/2 FQHE

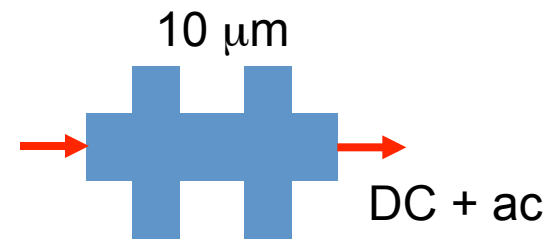
Down to 6.9 mK:

Tingxin Li

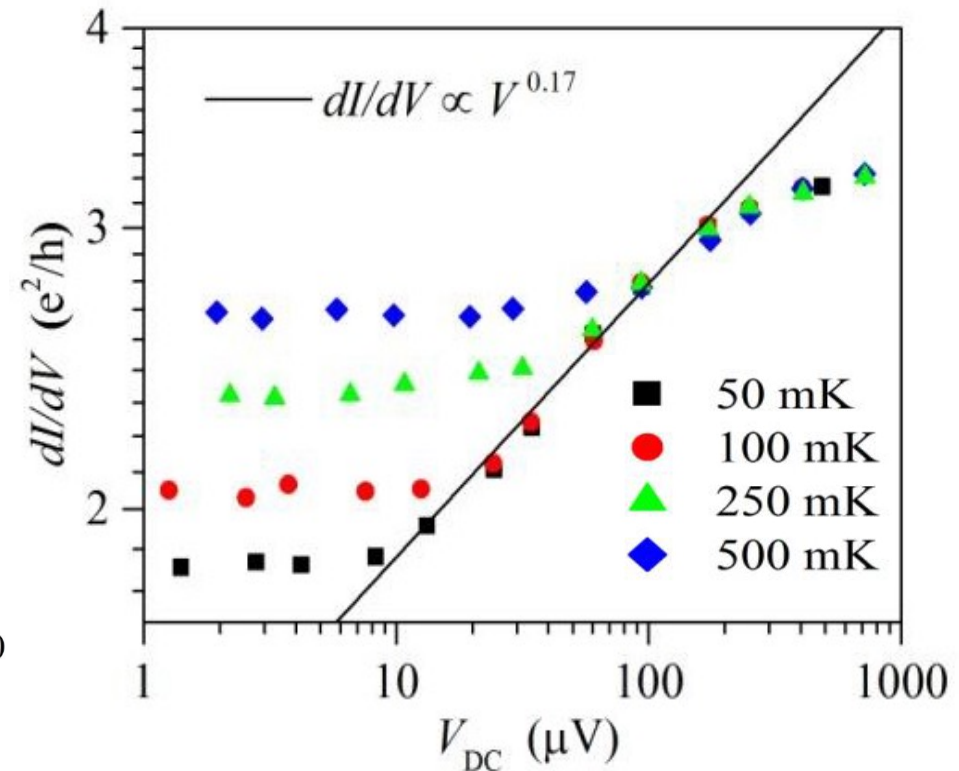
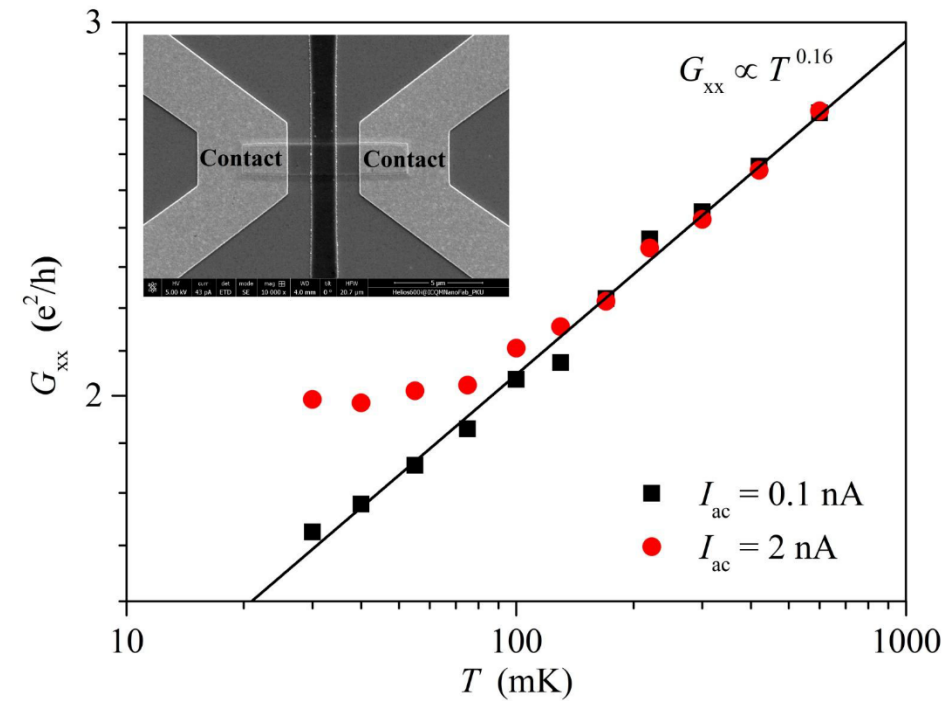


4-terminal Resistance: $L = 10 \mu m$
Strongly nonlinear with bias current

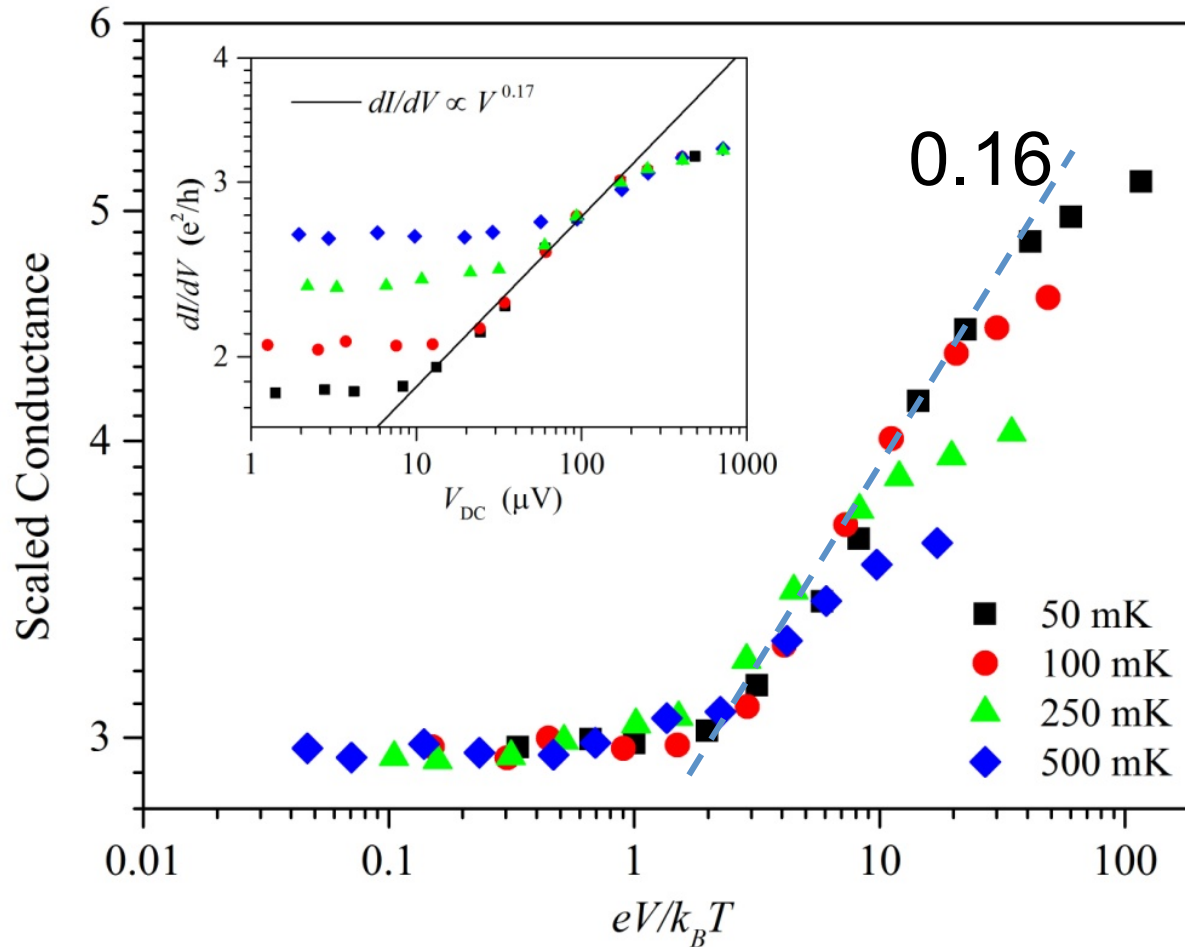
Voltage drops across $eV = R_{xx} * I_{ac}$



Temperature or DC bias -dependent

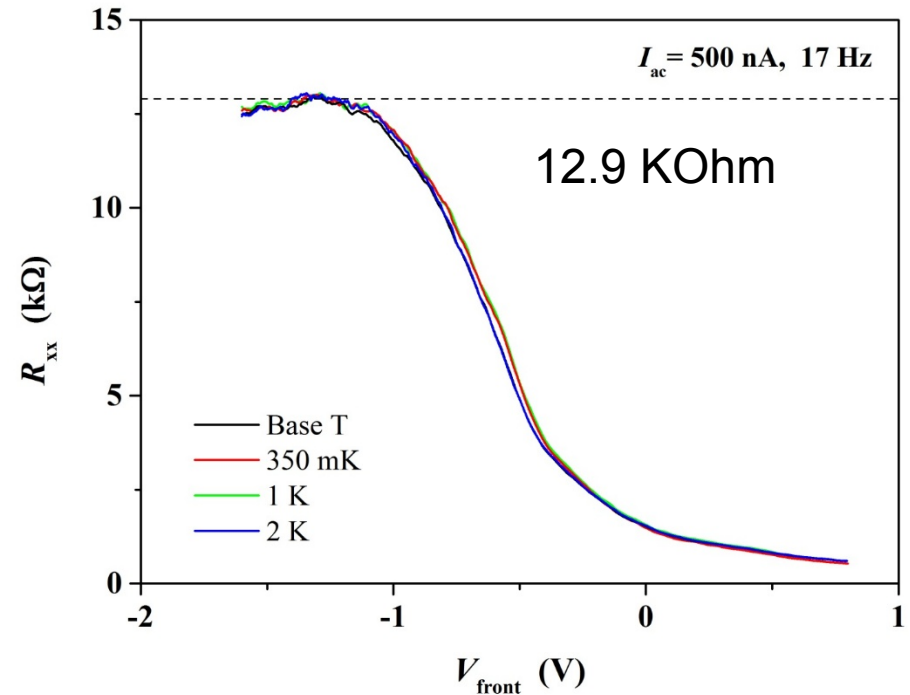
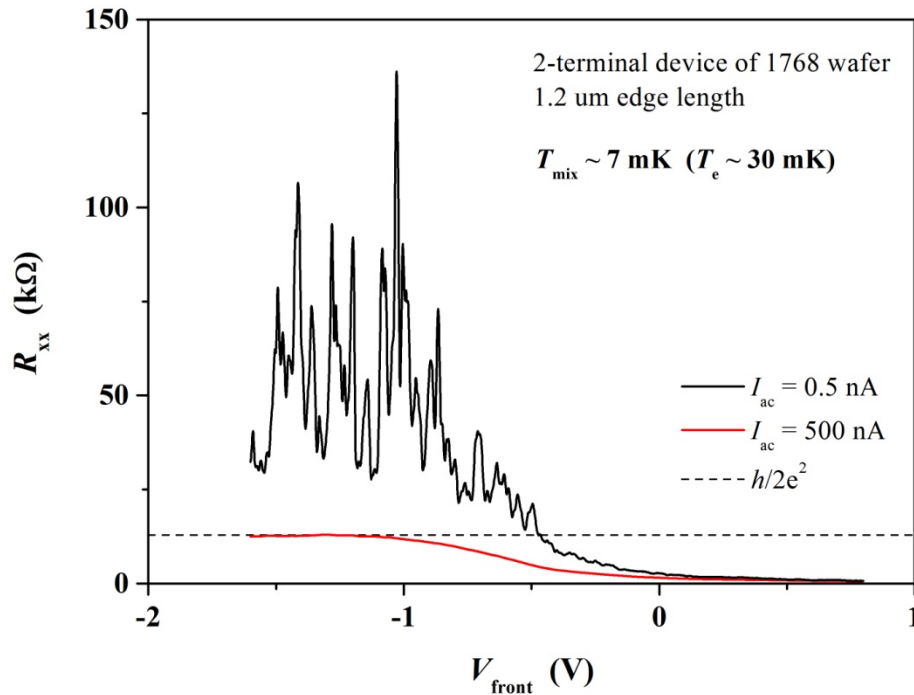


Wafer A Scaling Relation



- **Conductance and differential conductance scale as power-laws with temperature and bias voltage, respectively.**

Wafer B, quantized plateaus

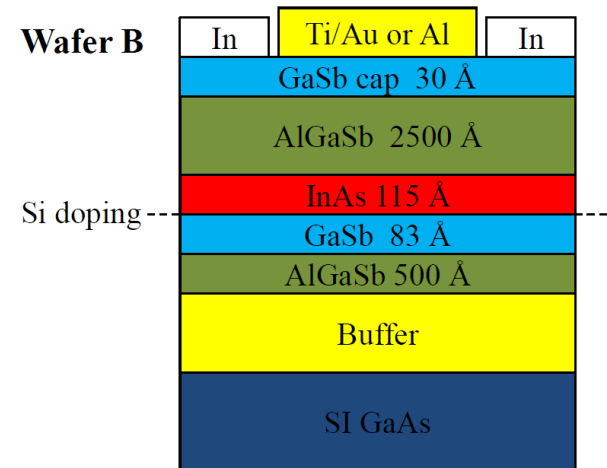


$L = 1.2 \mu\text{m}$, two terminal

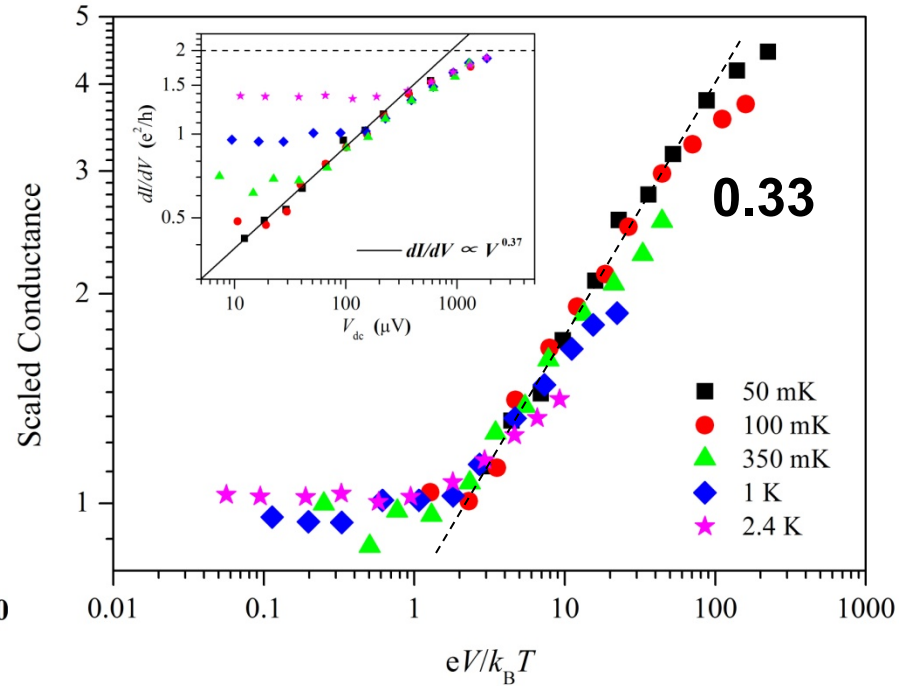
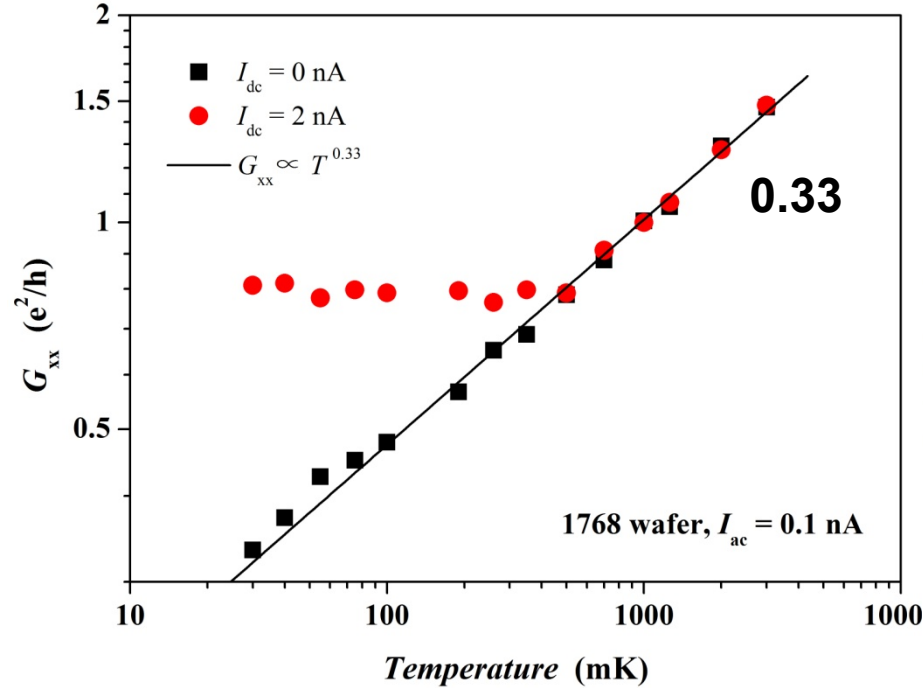


- There is no change for the quantized plateaus from base temperature to about 3K under large current.

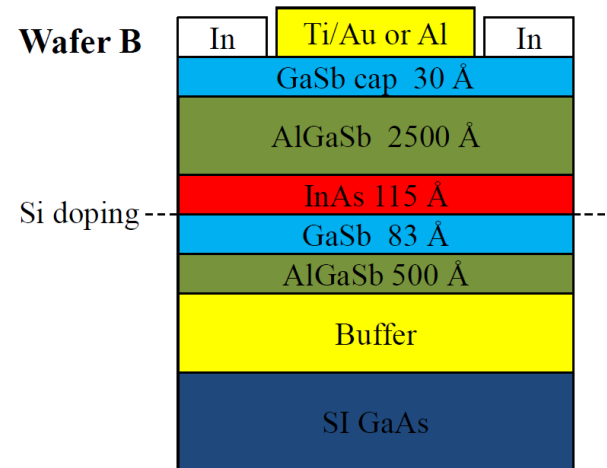
Wafer B



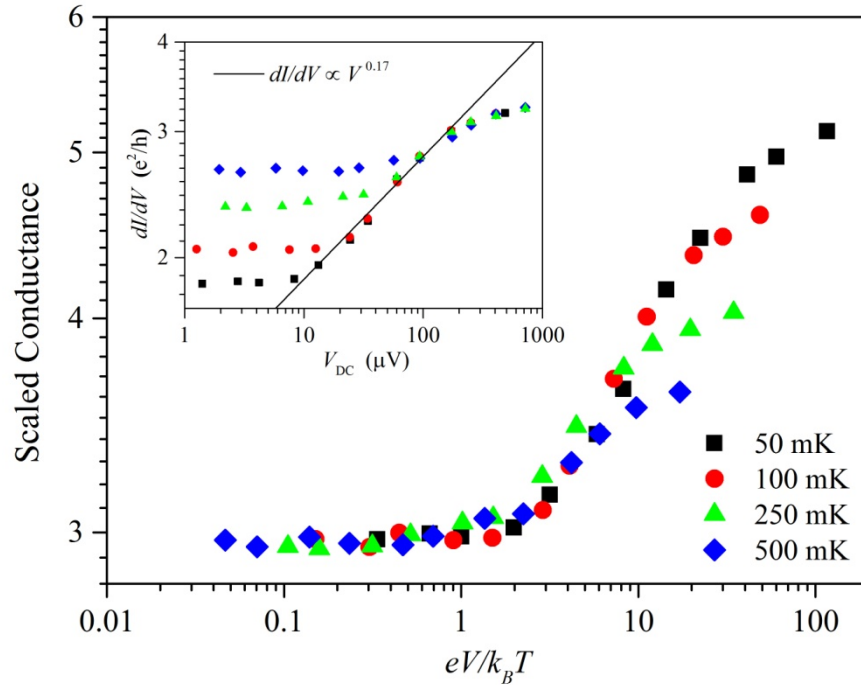
Wafer B, LL behavior



- Wafer B show stronger power-law behavior than wafer A.

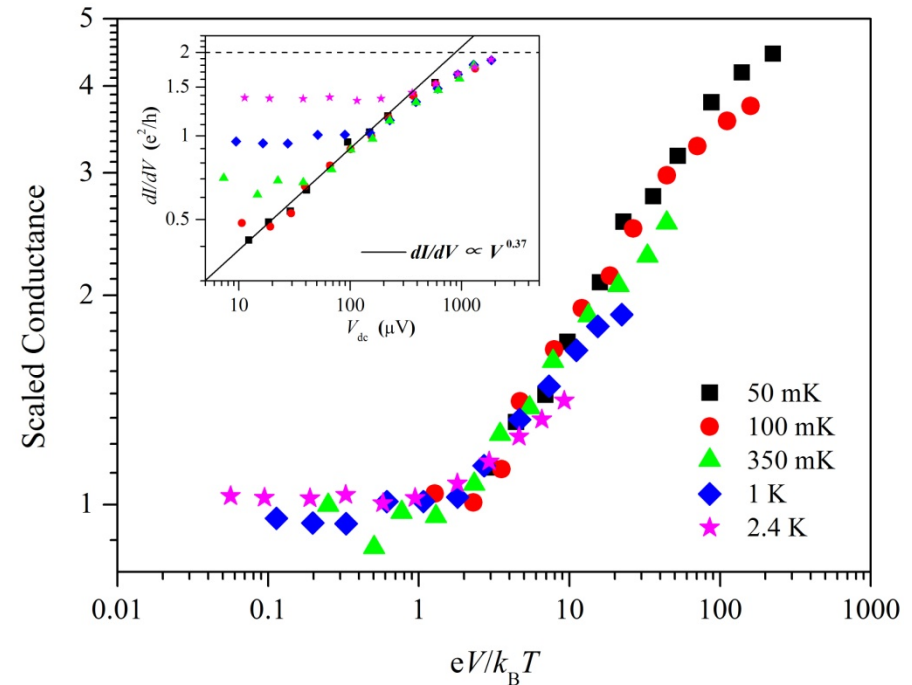


Comparison of two wafers



Wafer A Power exponent $\alpha \sim 0.16$,

$$\alpha = 2(1/4K-1) \rightarrow K \sim 0.23$$



Wafer B Power exponent $\alpha \sim 0.33$,

$$\alpha = 2(1/4K-1) \rightarrow K \sim 0.21$$

- Both two wafers tend to form a charge insulator at $T = 0$ and $V = 0$, consisting with the theoretical predictions for strong interacted helical edge states, i.e. Helical Luttinger-Liquid.

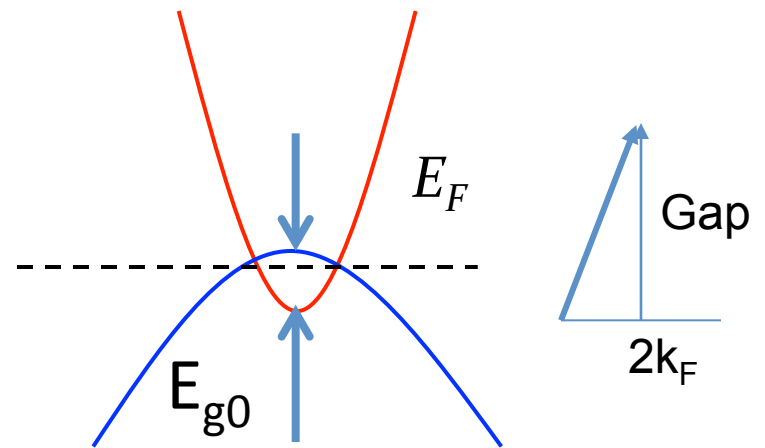
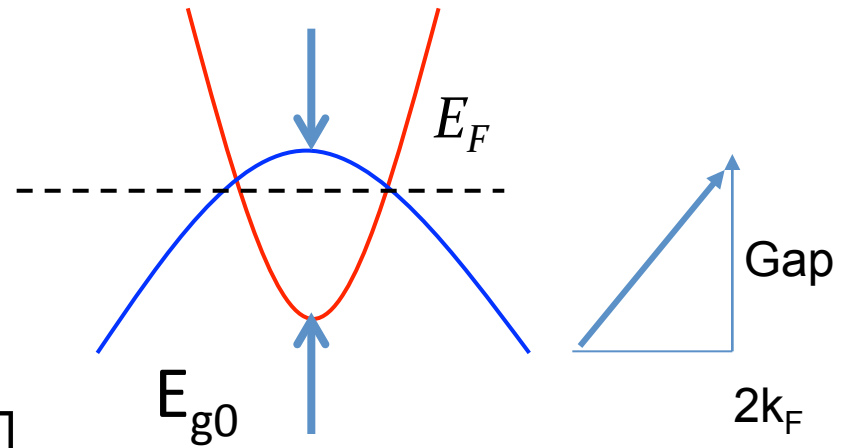
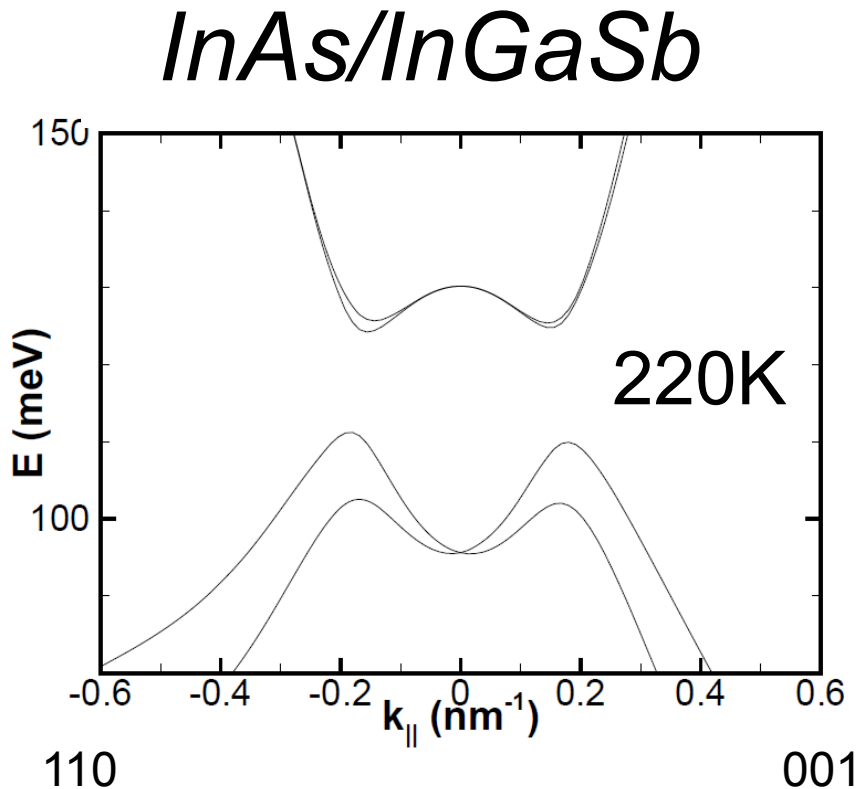
InAs/GaAs edge transport can be explained based on Luttinger Liquid

Microscopic details to be understood

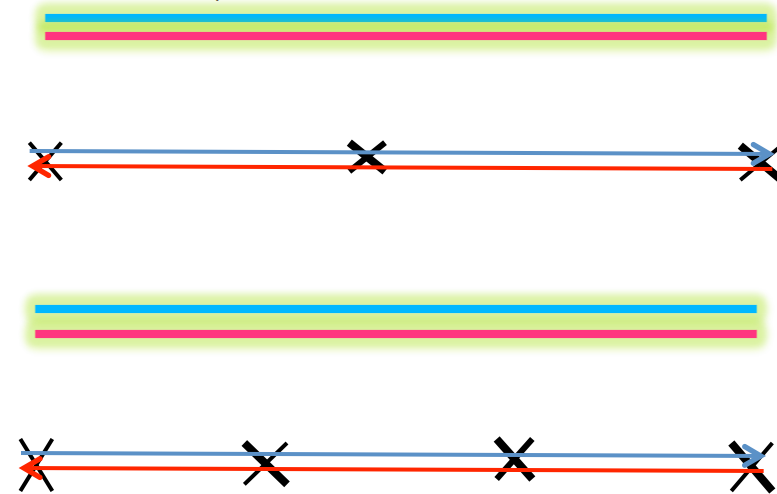
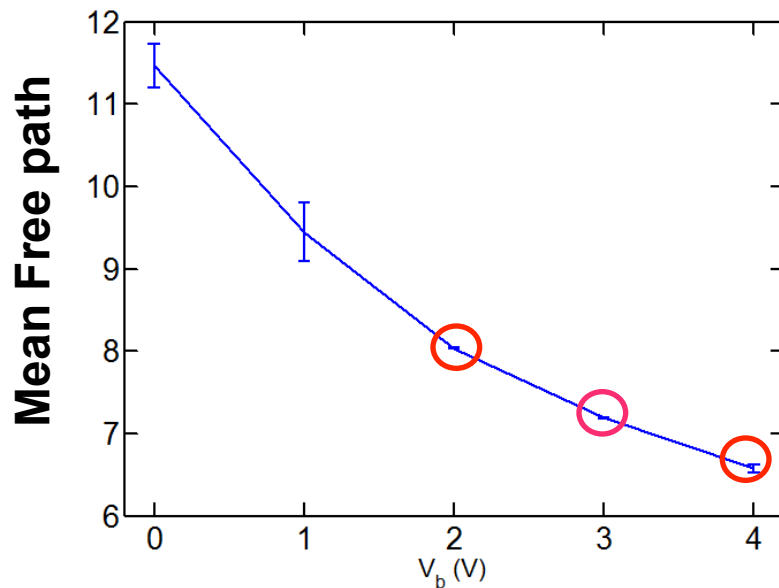
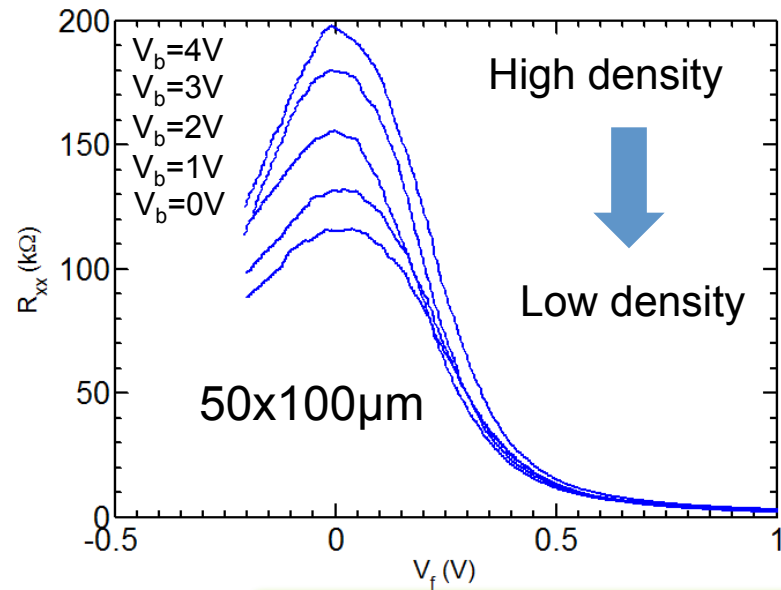
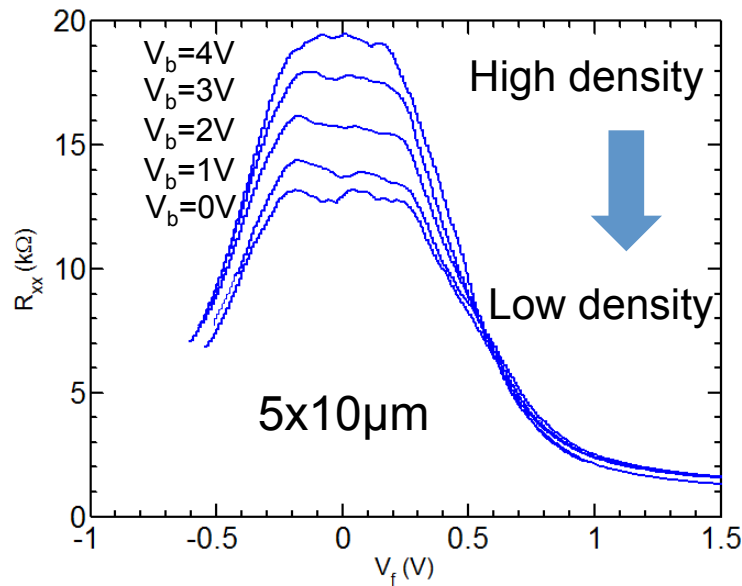
- 1) **Short edges** with Fermi liquid leads: Quantized conductivity $2e^2/h$
- 2) **Long edges** with Fermi liquid leads: reduced cond. from $2e^2/h$
- 3) $eV \gg k_B T$ Large bias regime: no T -dependence
 $eV \ll k_B T$ Small bias regime: no eV -dependence
 $eV \sim k_B T$ Conductance scaling in $eV/k_B T$

Outlook: Tuning the Interactions

Strain Layer Engineer. + Gating

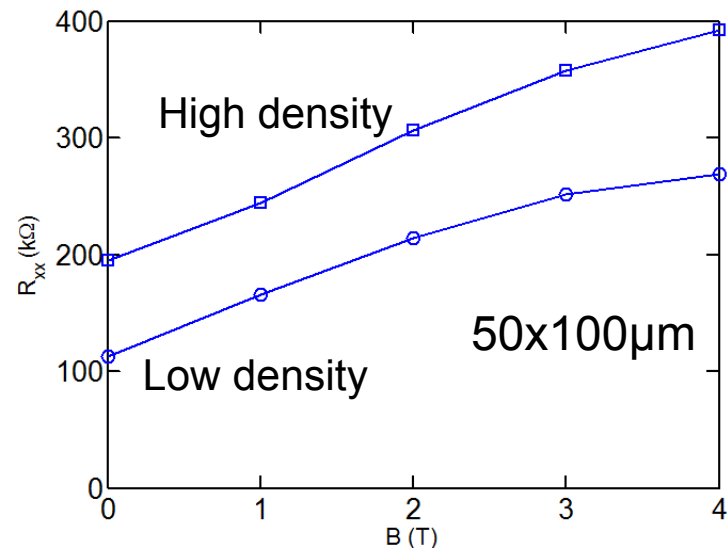
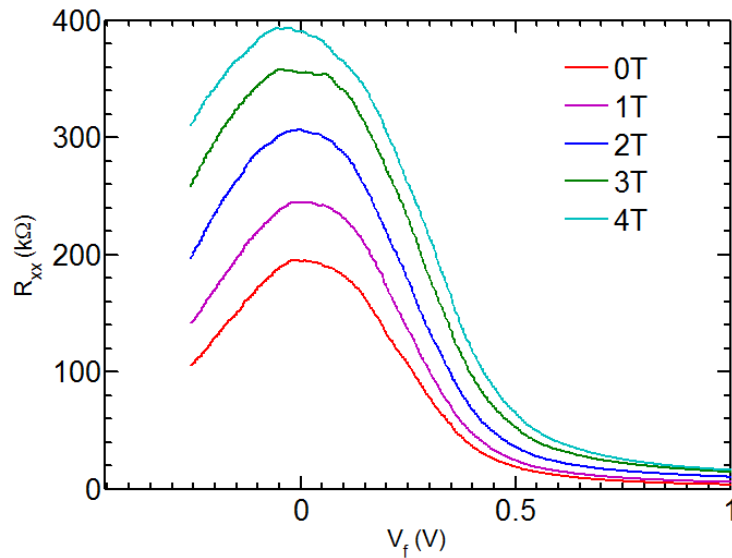


“Mean Free Path” Increases with V_F

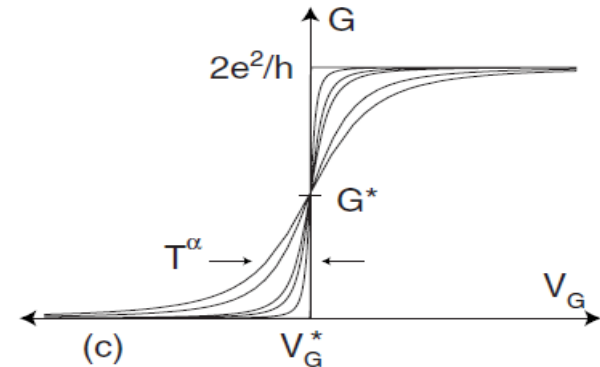
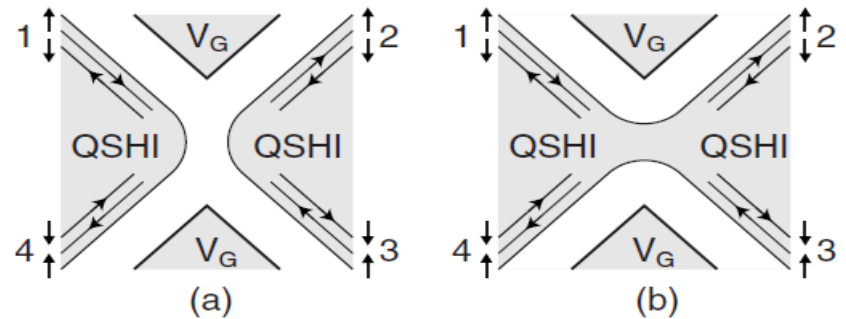
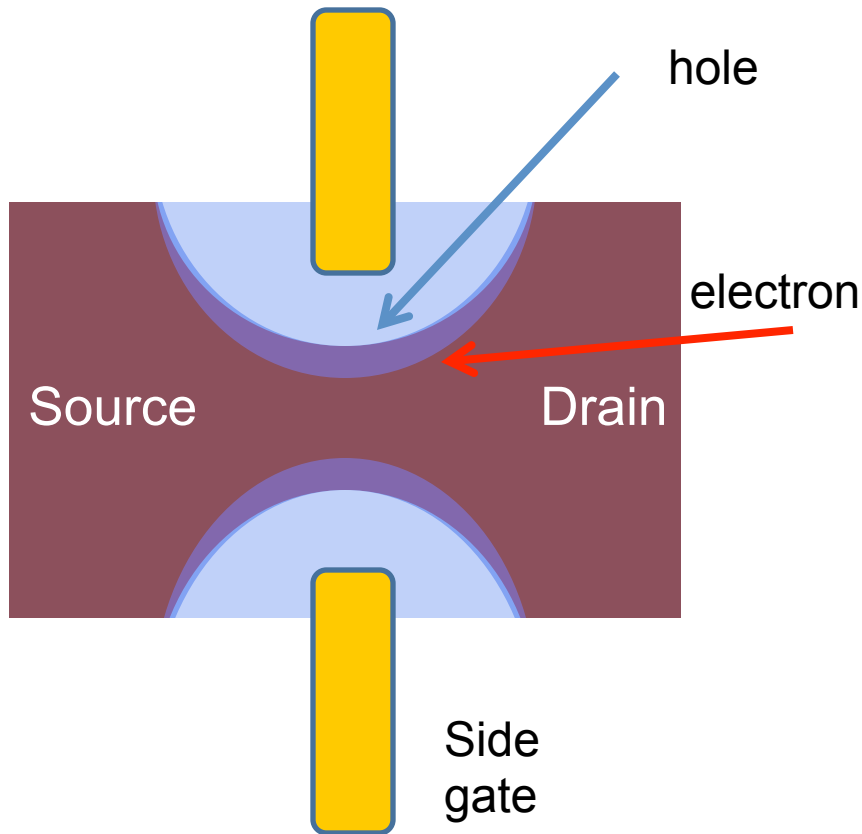


Magnetic Response Depends on V_F

Less Interacting Edge States

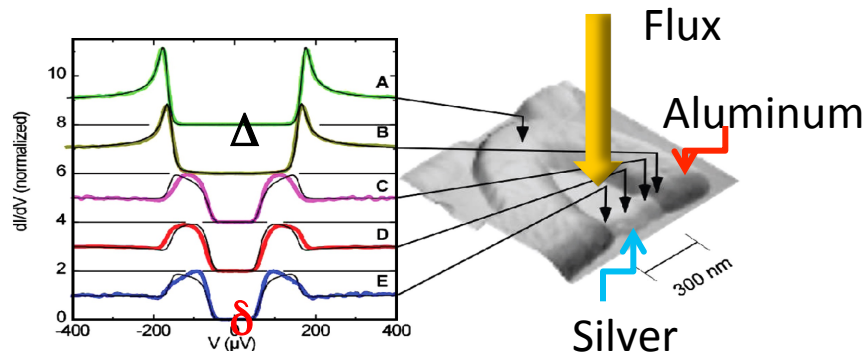


Quantum Point Contact

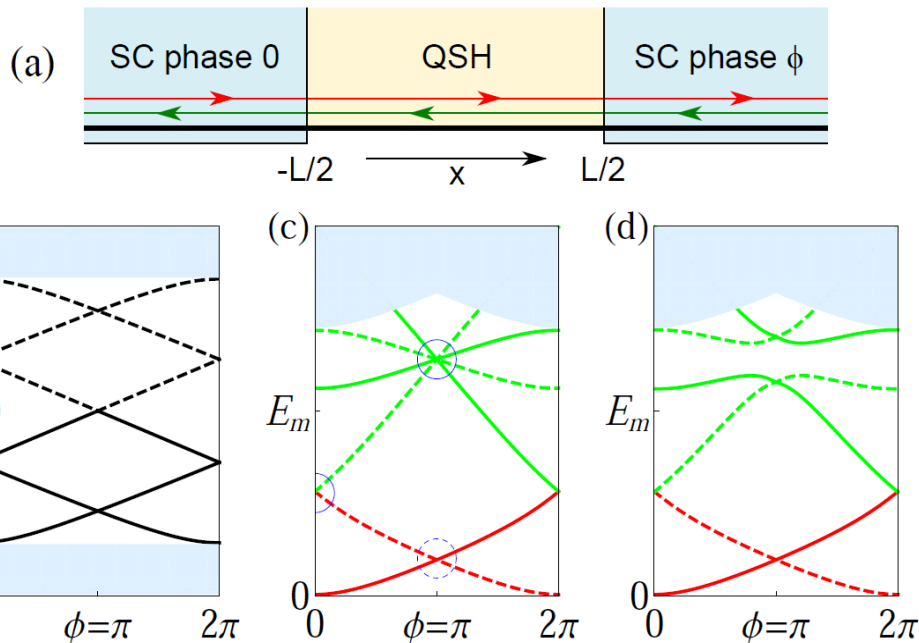


Shot Noise $e/2$
Quantum Critical Point

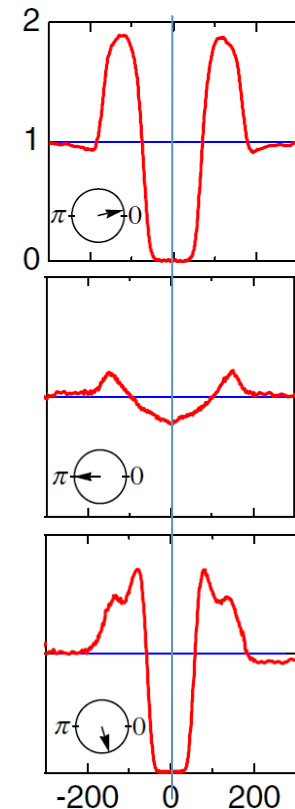
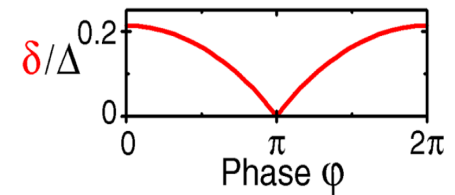
Tunneling Into Phase Controlled Junction



H. le Sueur et al, PRL 08



Phase Control



$8-\pi i$
Josephson
Effect due to
Interaction

Zhang & Kane
PRL 13

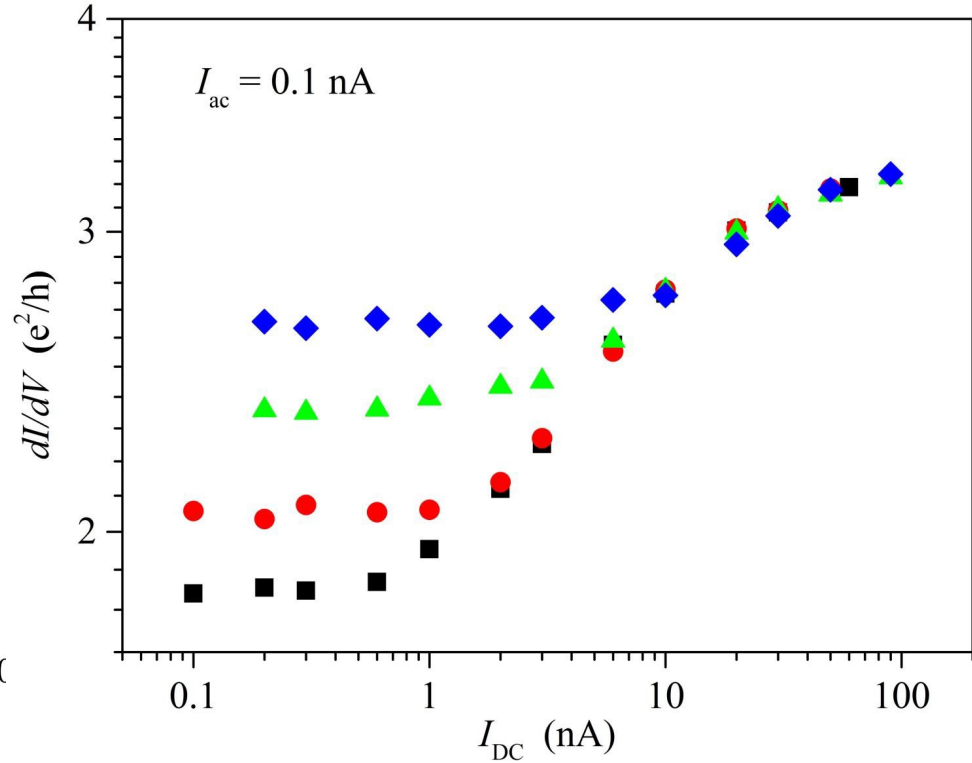
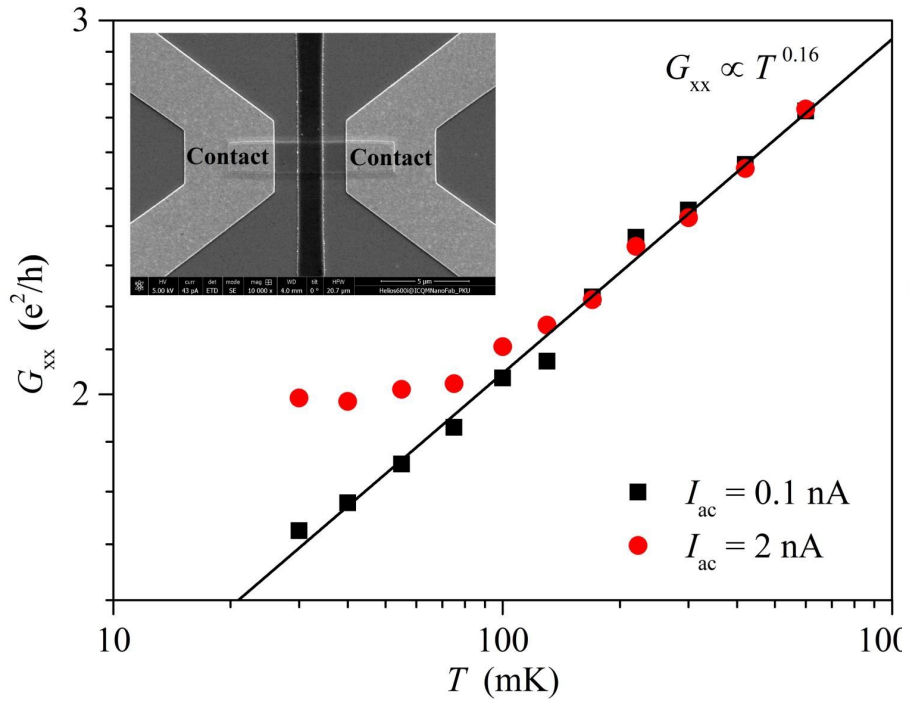
SUMMARY

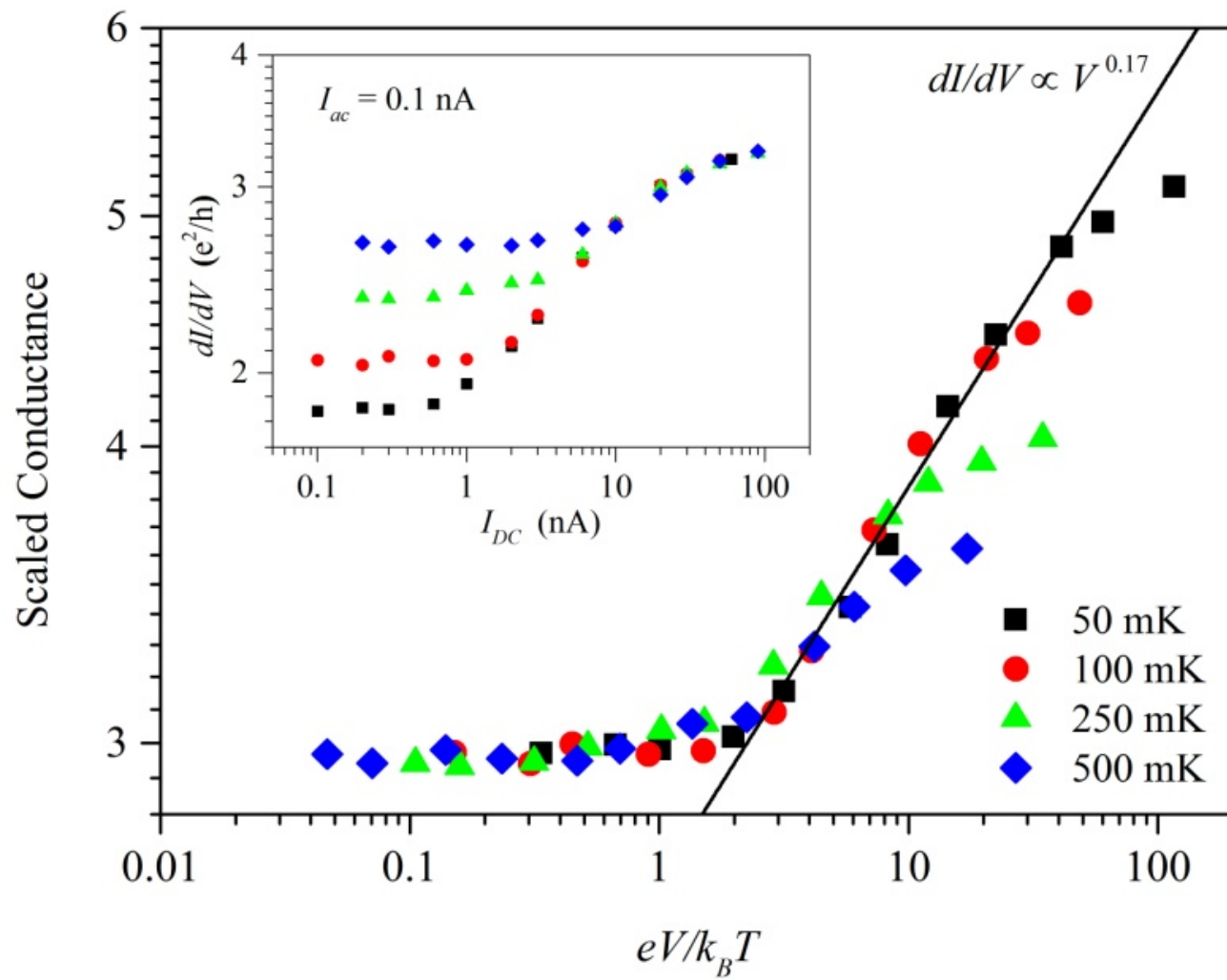
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- 3) Strong Evidences for Helical Luttinger Liquid
Microscopic details to be understood

New opportunities to study 1D correlation effects and related fractional charge and fractional Josephson effect

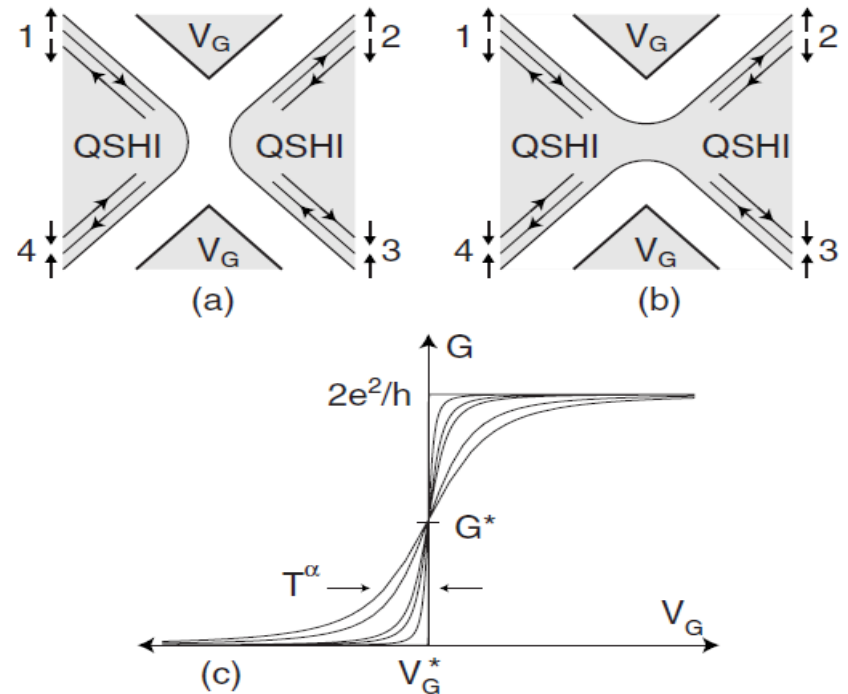
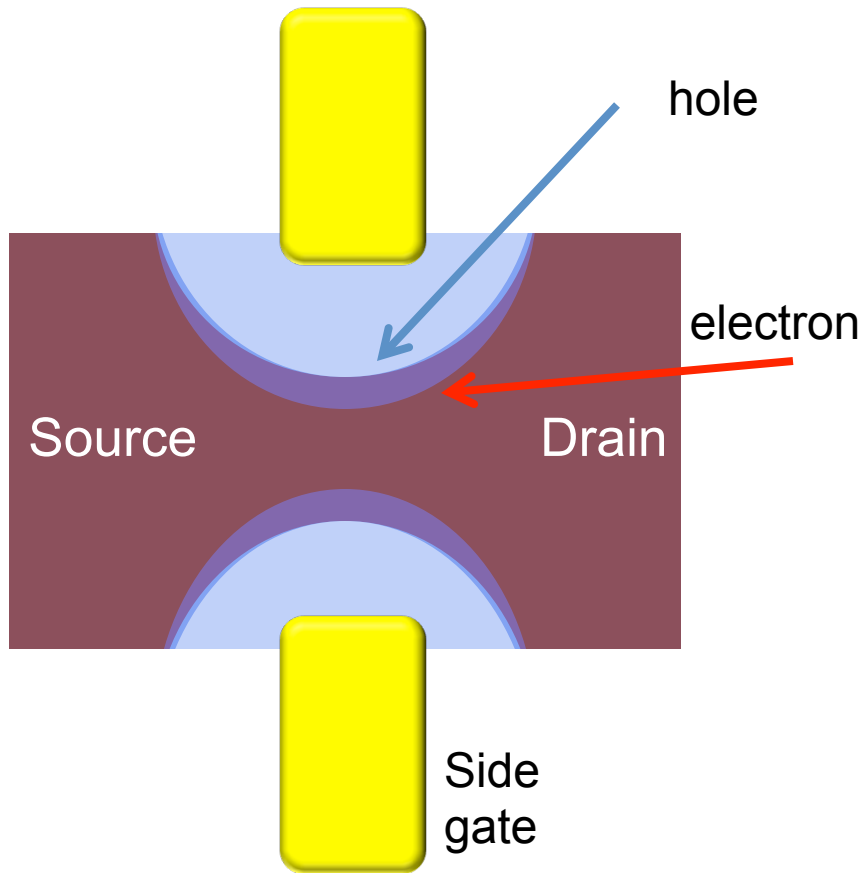
THANKS

Temperature-dependent in small-bias regime



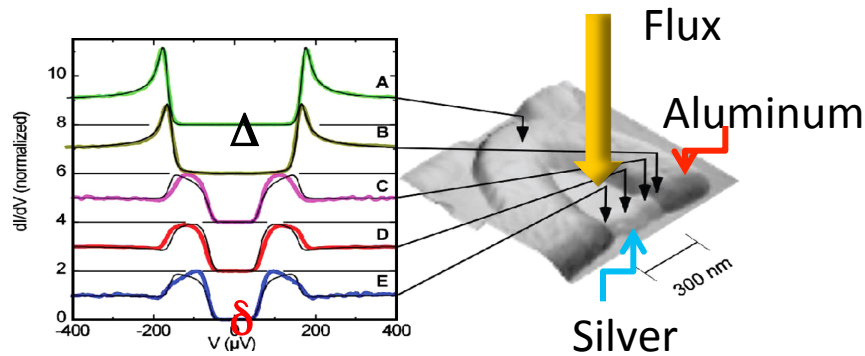


Quantum Point Contact

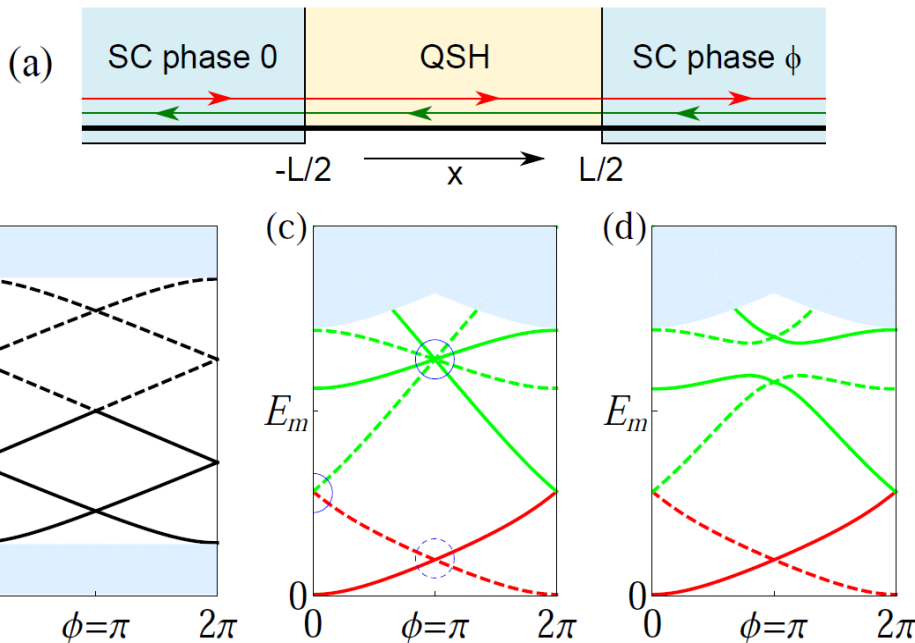


Shot Noise $e/2$
Quantum Critical Point

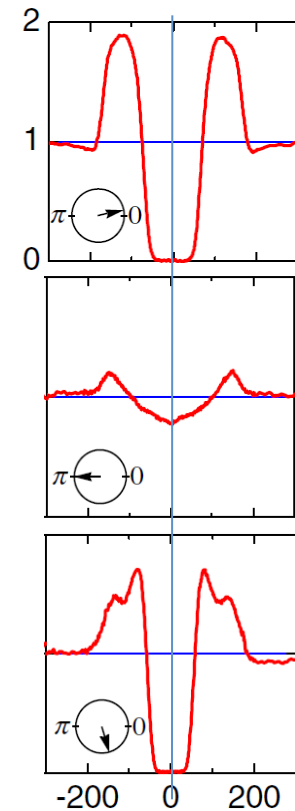
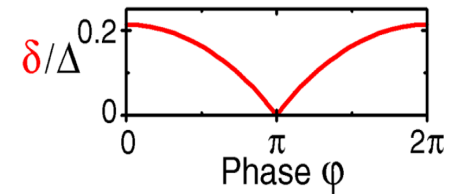
Tunneling Into Phase Controlled Junction



H. le Sueur et al, PRL 08



Phase Control



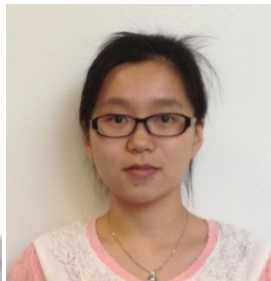
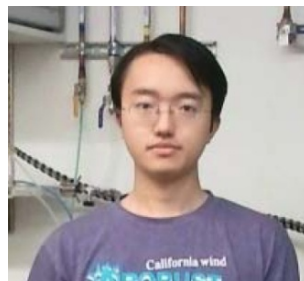
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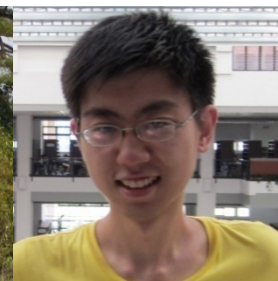
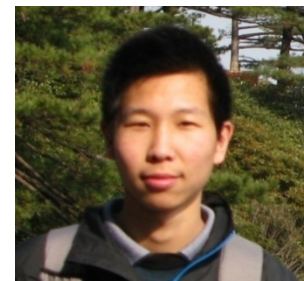
Rice

Ivan Knez Lingjie Du Ruiyuan Liu Jie Zhang



Peking U

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Collaborations

Teledyne

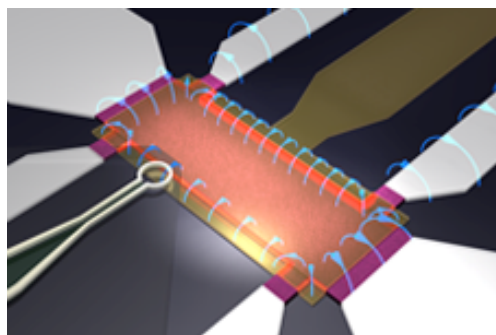
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Among many others

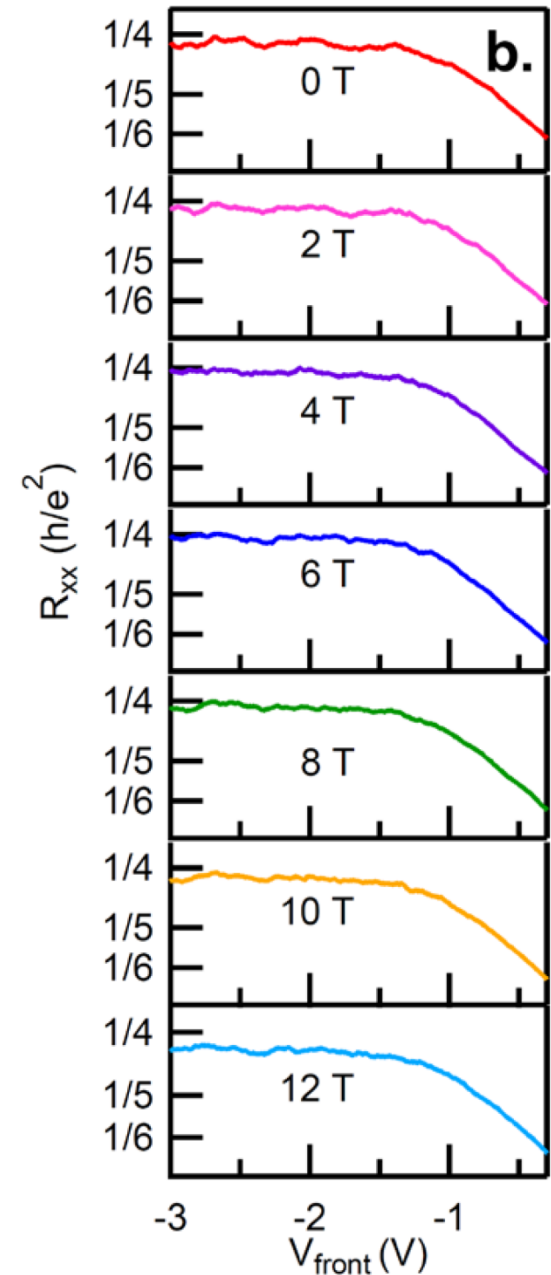
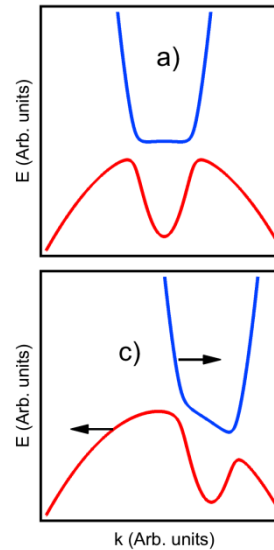
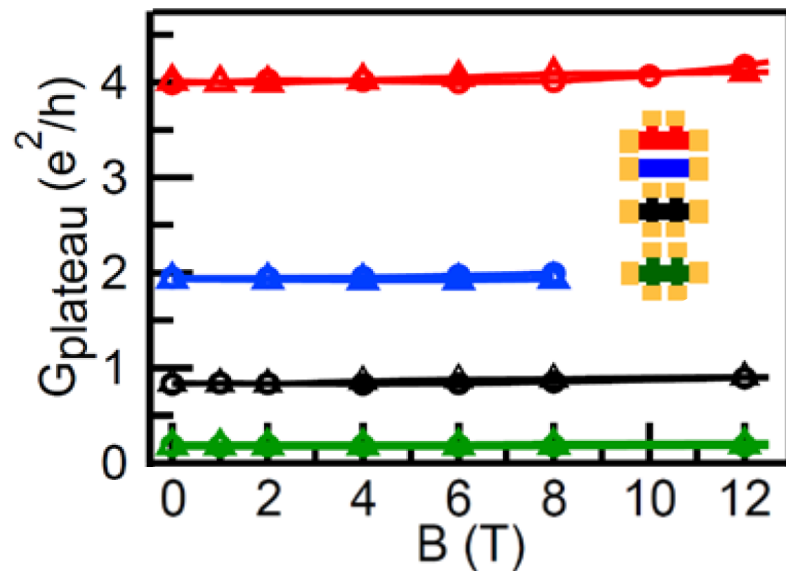
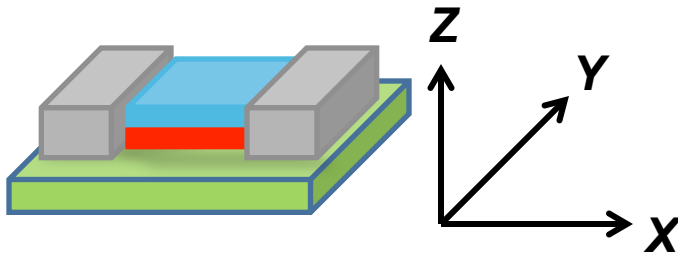
SUMMARY

- 1) Robust Quantum Spin Hall Effect
- 2) InAs/GaSb Helical Edge: clean and highly-tunable 1D system
- 3) Strong Evidences for Helical Luttinger Liquid must understood before engineering Majorana bound states
New opportunities to study 1D correlation effects and related fractional charge and fractional Josephson effect

THANKS

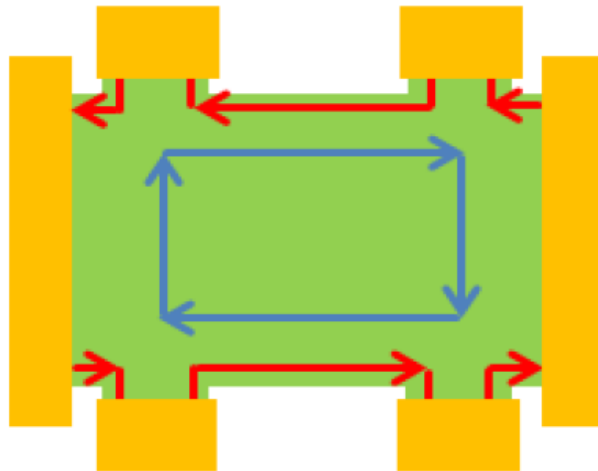
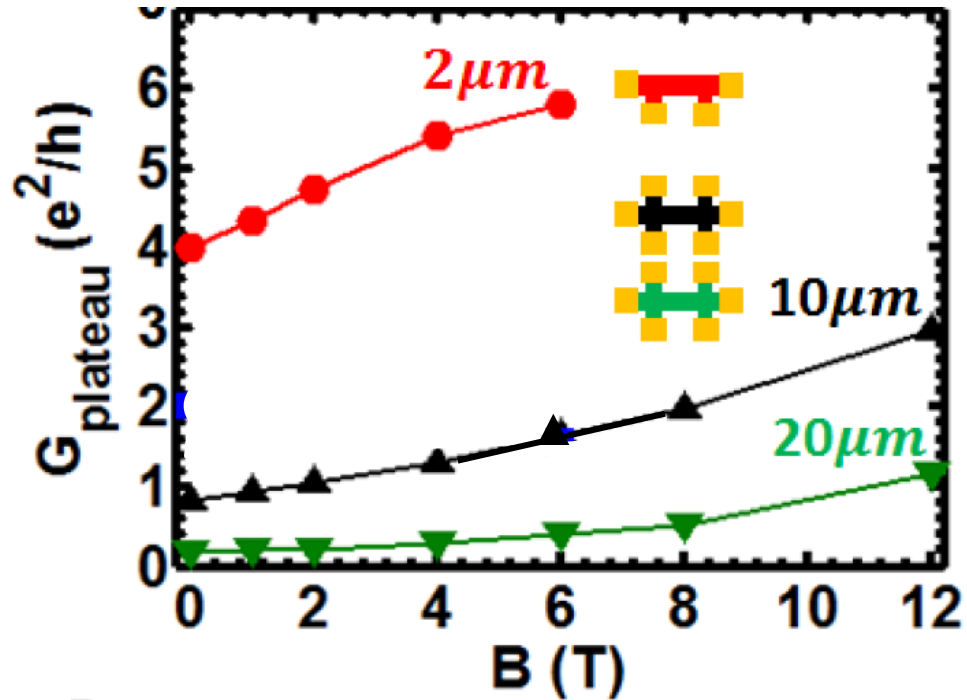
THANKS

In-Plane Magnetic Fields

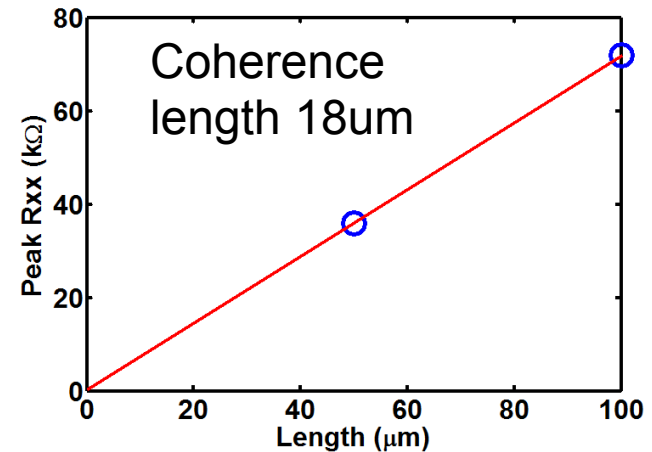
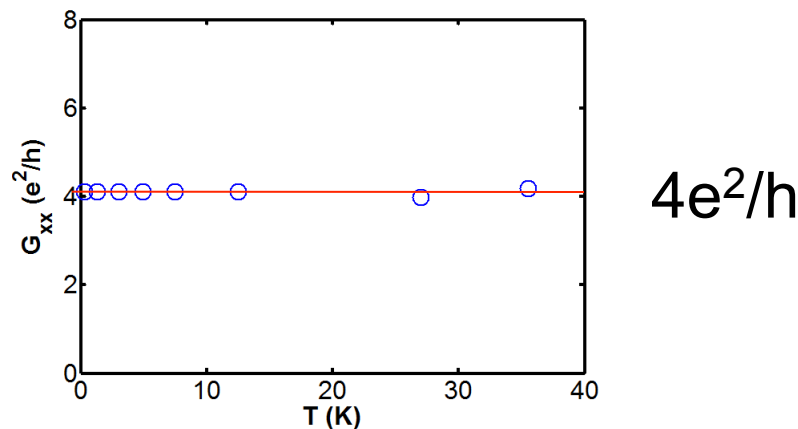
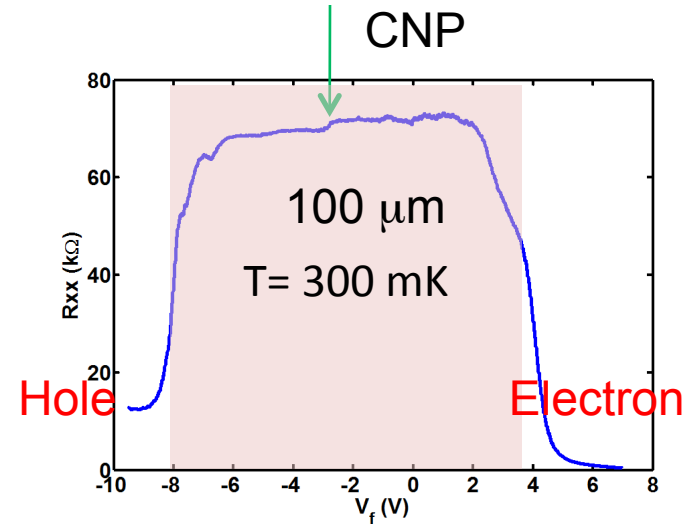
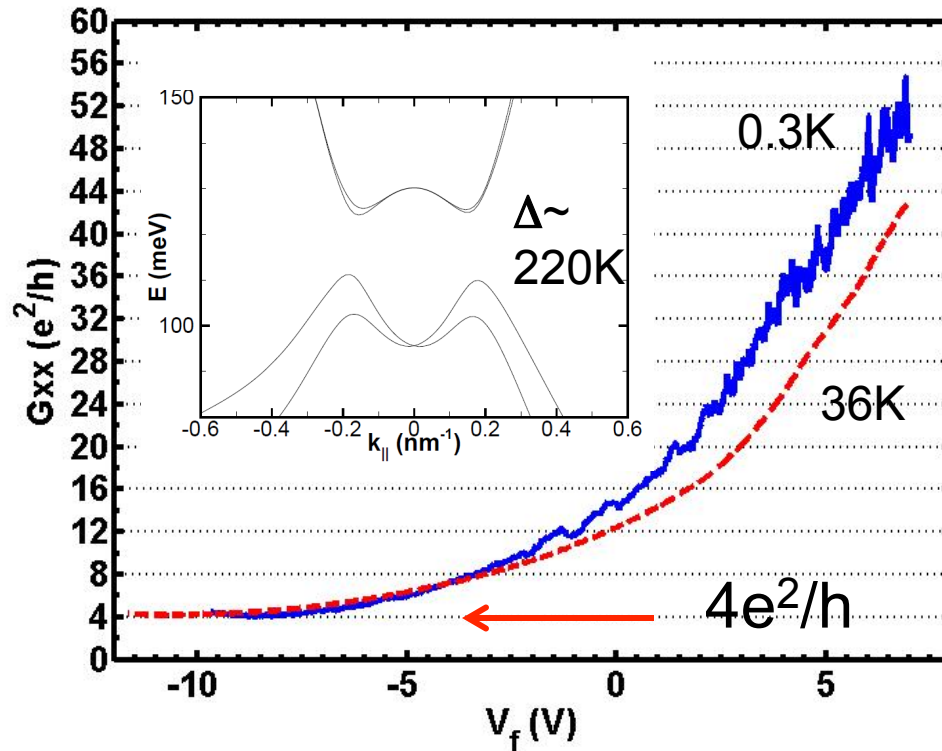


Small $g^* \sim 0.5$ PRB83, 155412 (2011)

Perpendicular Magnetic Field



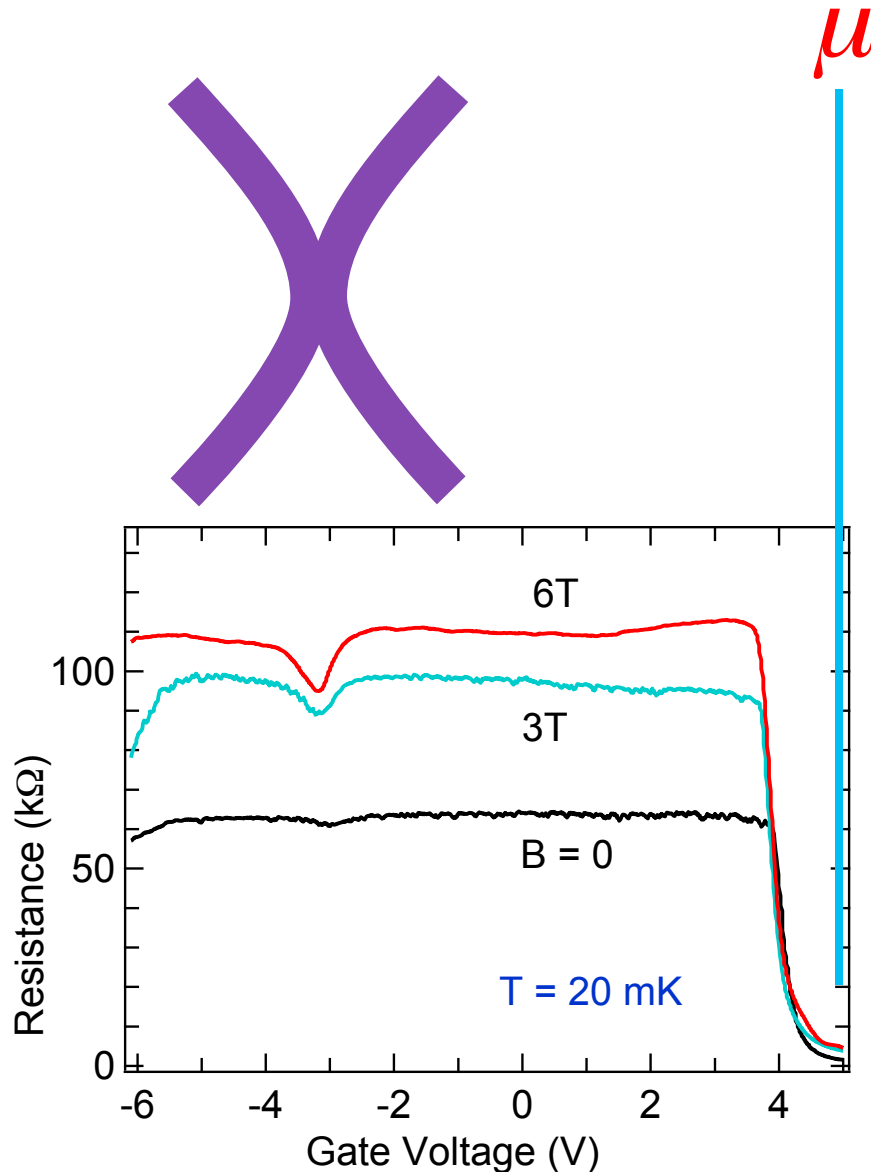
New Wafer (Strained QW) Preliminary Data



Quantization of edge in small sample up to 36K

$$\tau = 18\text{ }\mu\text{m} / 1 \times 10^5\text{ m/s} \sim 180\text{ ps}$$

Opening of Edge Gap by Perpendicular B



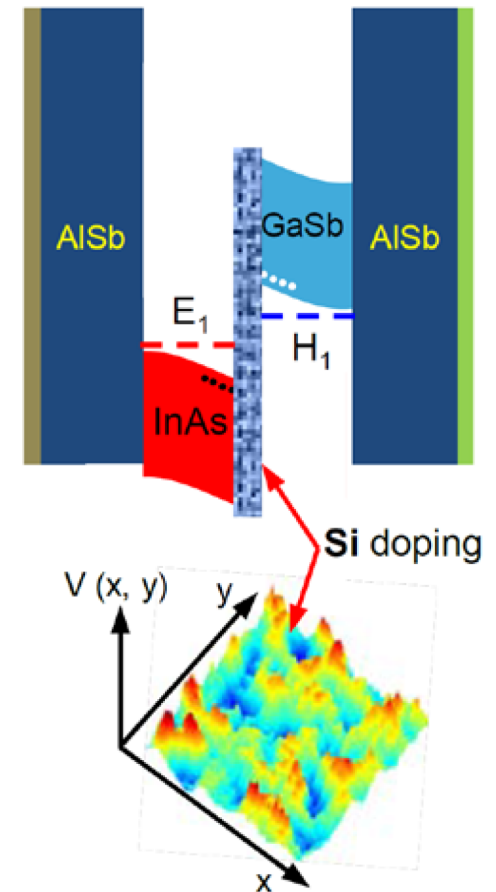
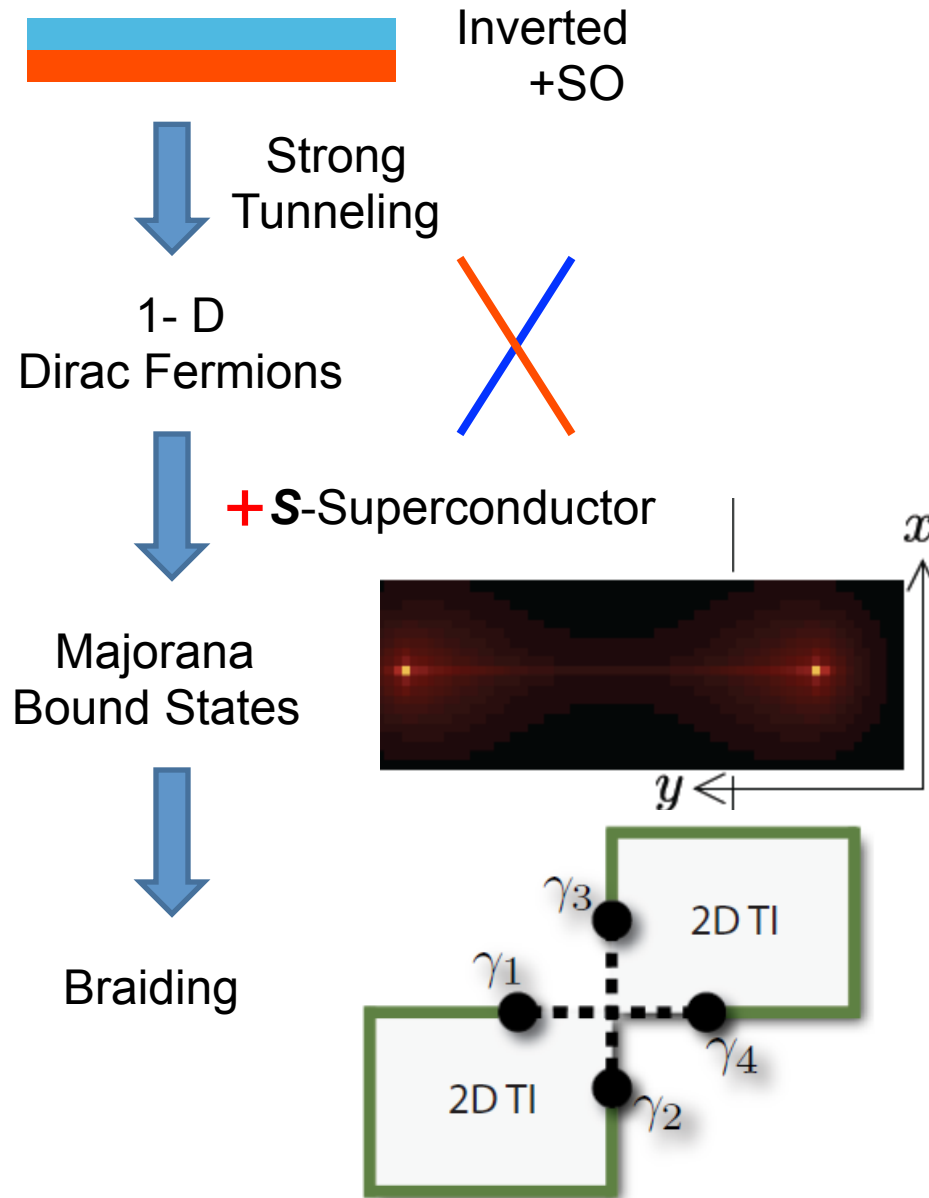
Strained QW of 100 μm length

- B = 0: Weak resistance dip already at Dirac point
- Dip deepens with B
- Gap not fully open

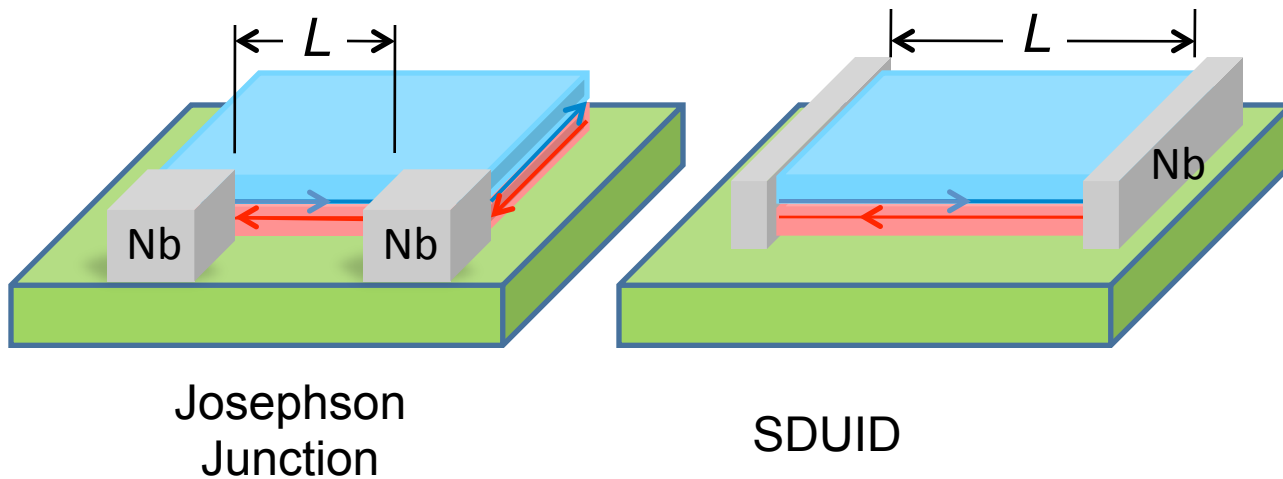
$$\Gamma \sim 1K$$

Tunneling DOS: STM

InAs/GaSb Majorana Platform



General Design



$$L = 0.1 \sim 1 \mu\text{m}$$

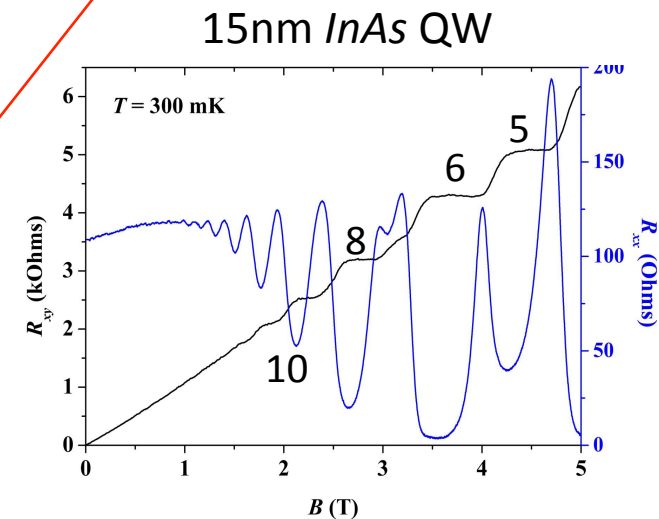
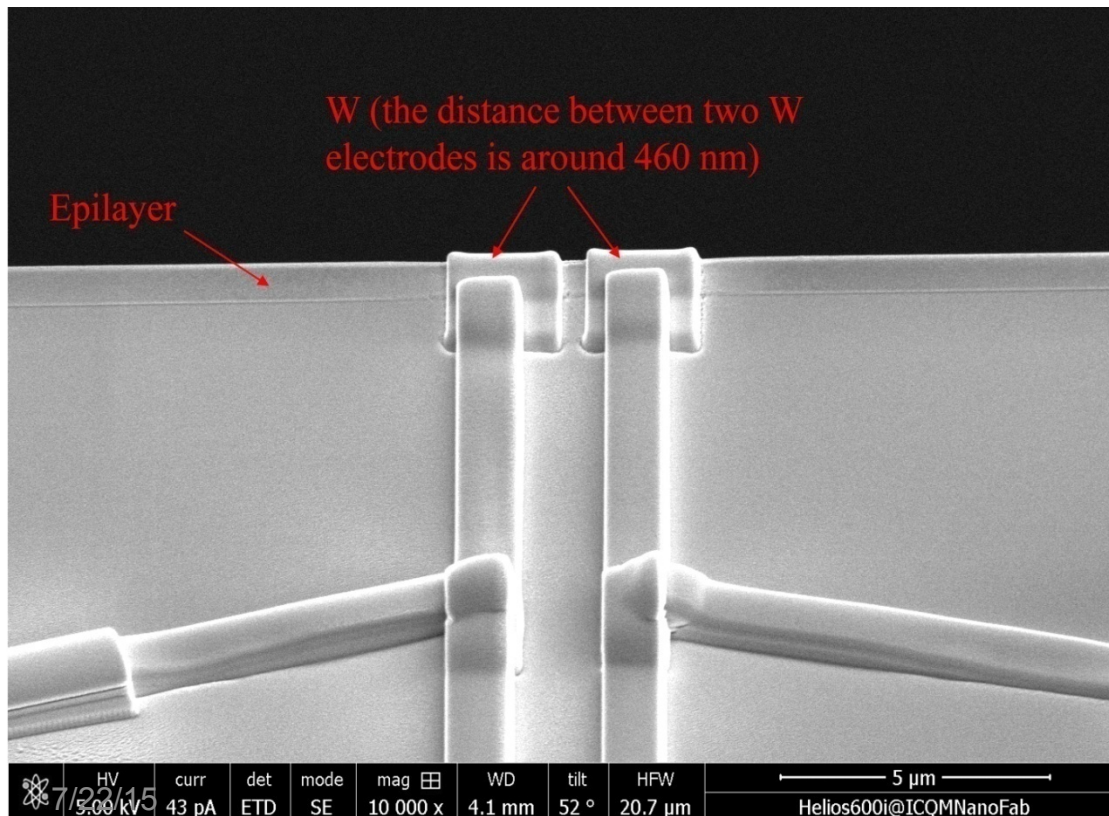
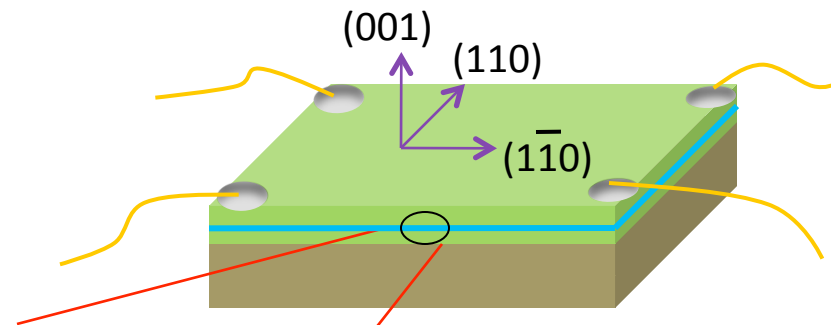
Direct Writing on Cleaved Edge

FIB-induced deposition of
superconducting nanowires

W(40%):C(40%):Ga(20%)

$T_c = 5.2\text{K}$, $H_c = 9.5\text{T}$

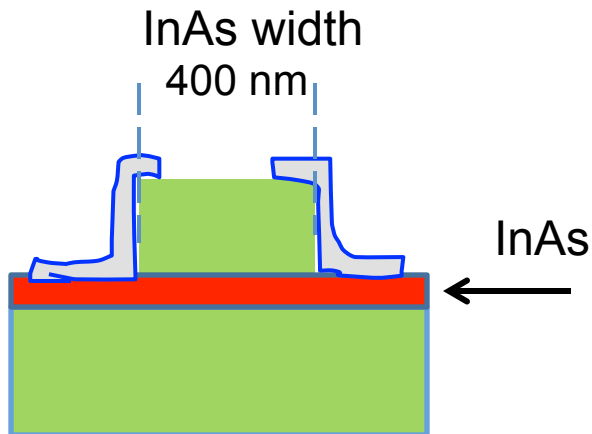
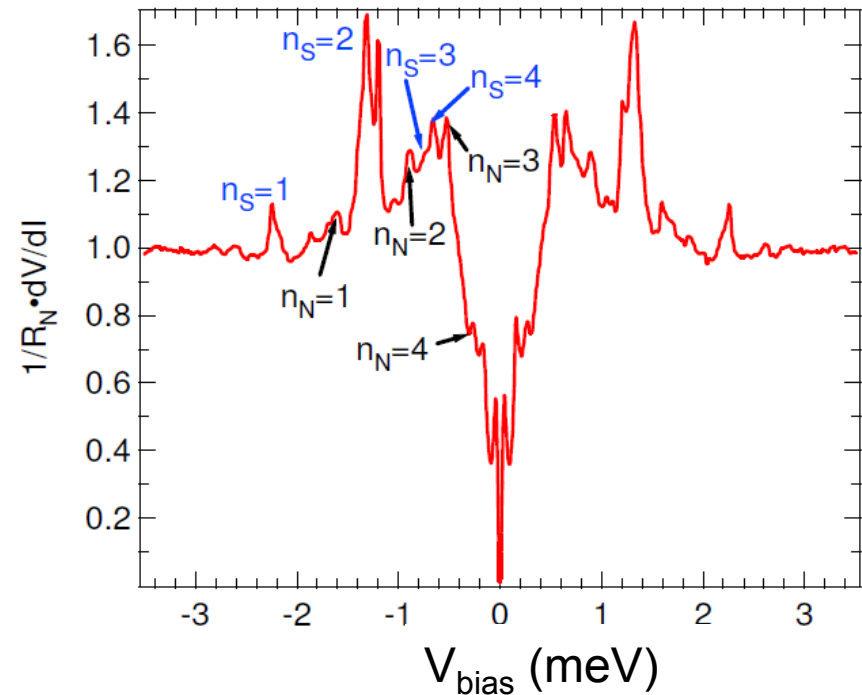
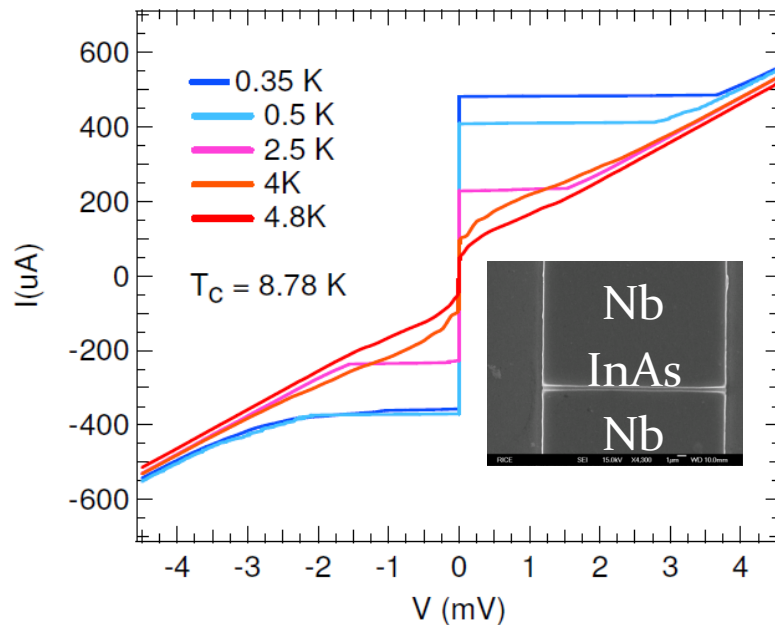
Sadki, Ooi, and Hirata, *Appl. Phys. Lett.* 85, 6206(2004)



Proximity effect
IQHE edge state

InAs is SC-Friendly: *Standard Nb-InAs-Nb*

Knez, Du, Sullivan 09



Multiple Andreev Reflection:

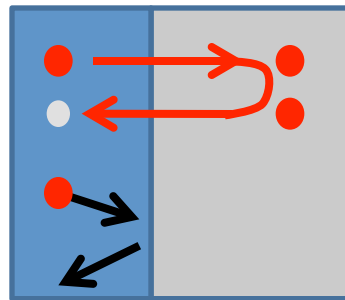
3 sets of peaks: Δ_N , Δ_S , $\Delta_N + \Delta_S$

$$V_{SGS} = \frac{2 \cdot \Delta_S}{n \cdot e}, \frac{2 \cdot \Delta_N}{n \cdot e}, \frac{2 \cdot (\Delta_S \pm \Delta_N)}{n \cdot e}$$

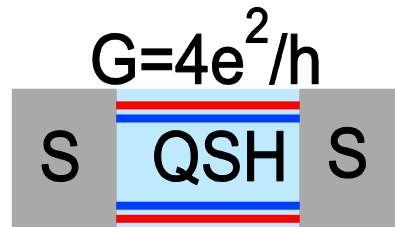
$$\Delta_N / \Delta_S \sim 0.67$$

$$\Delta_S = 1.31 \text{ meV}; \Delta_N = 0.87 \text{ meV}$$

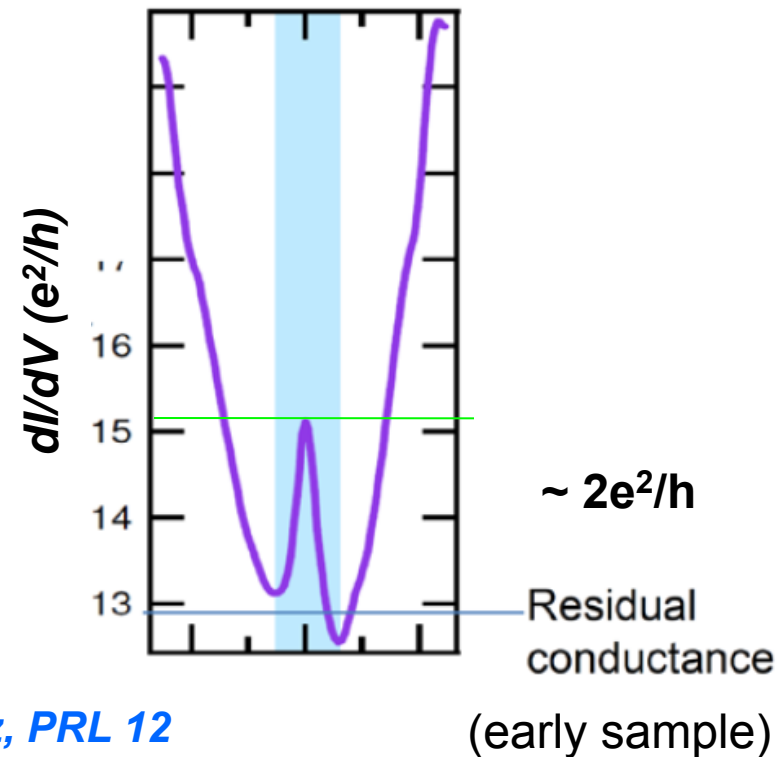
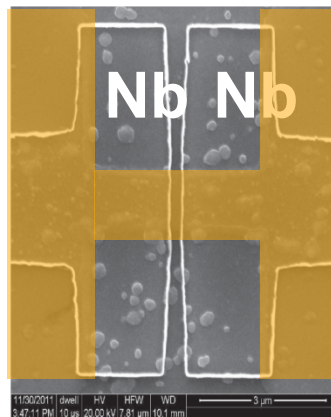
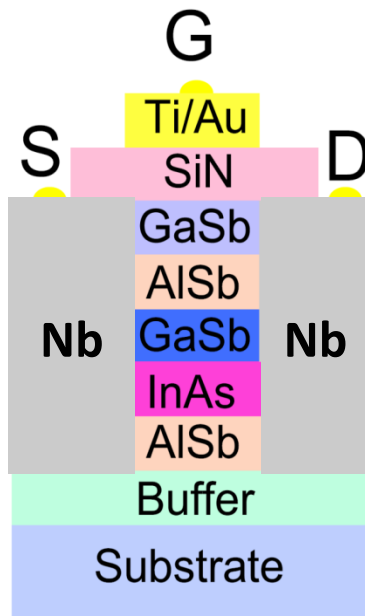
Andreev Reflection of Nb-(InAs/GaSb)-Nb Junction



Normal SC



$$G = 2e^2/h$$



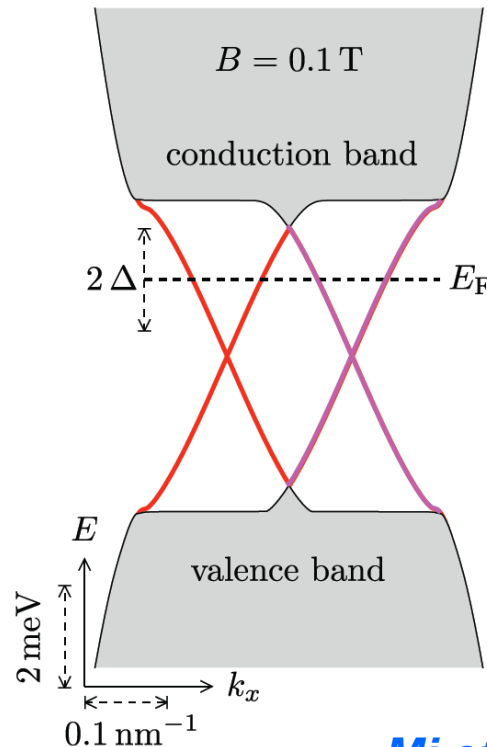
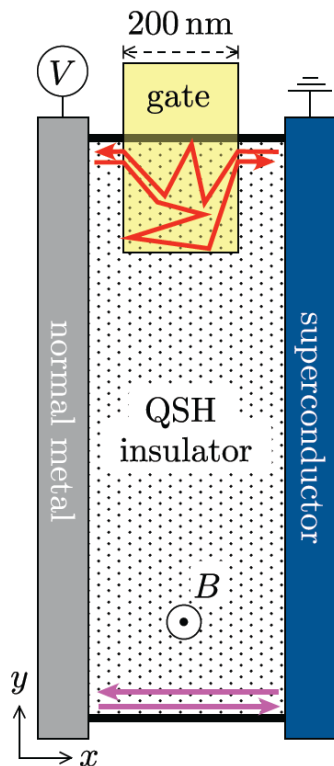
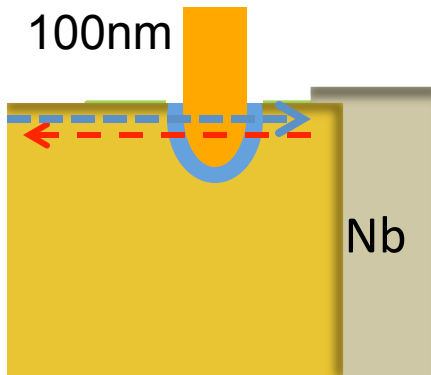
Knez, PRL 12

Goal: Creating MBS on the Edge

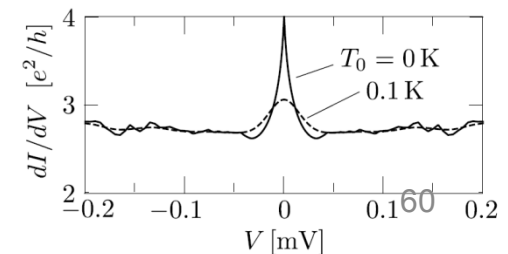
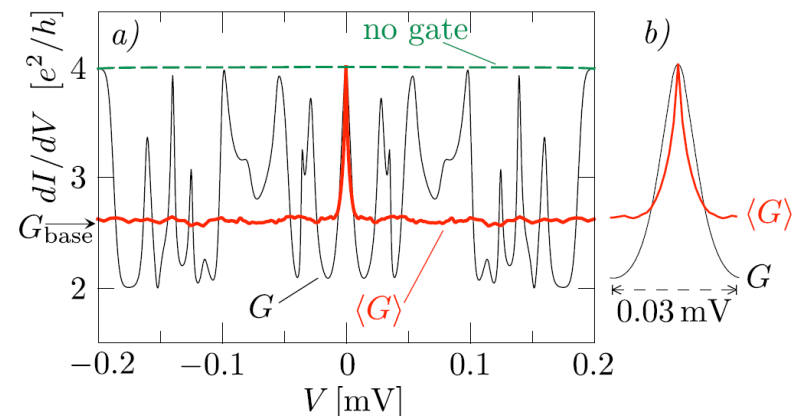
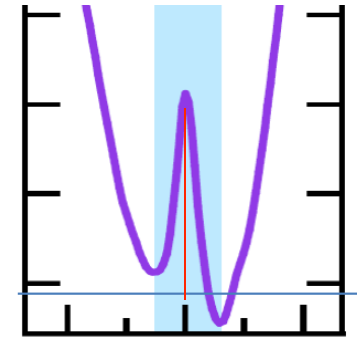
Main challenges:

1. Supercurrent -- small v_F
2. Soft induced-SC gap -- interface roughness etc
2. Terminating edge --

On-Going Experiment towards Isolation

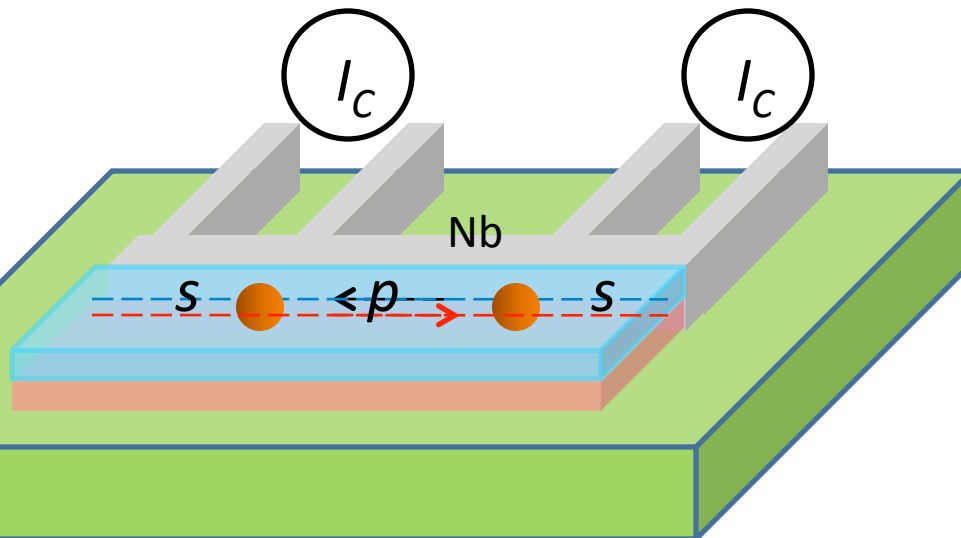
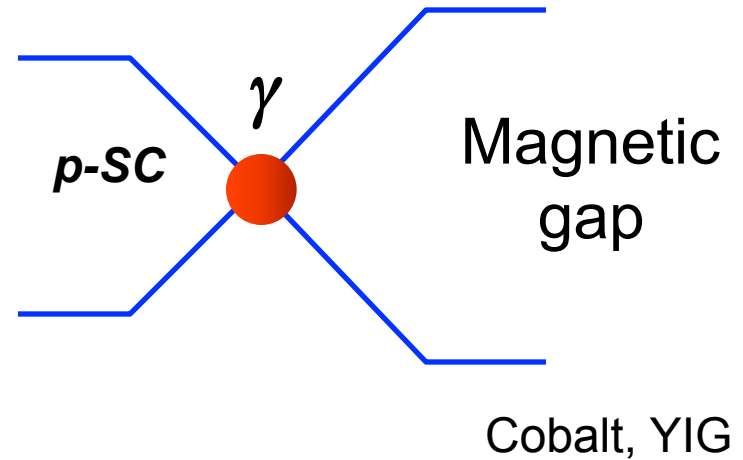
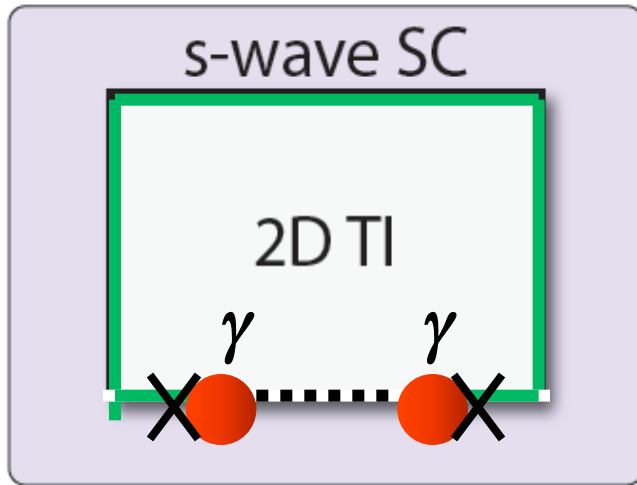


1. MF delocalize
2. Trapping MF by metallic dot + B
3. Backscattering confines MF
4. Zero modes stands out in disorder average



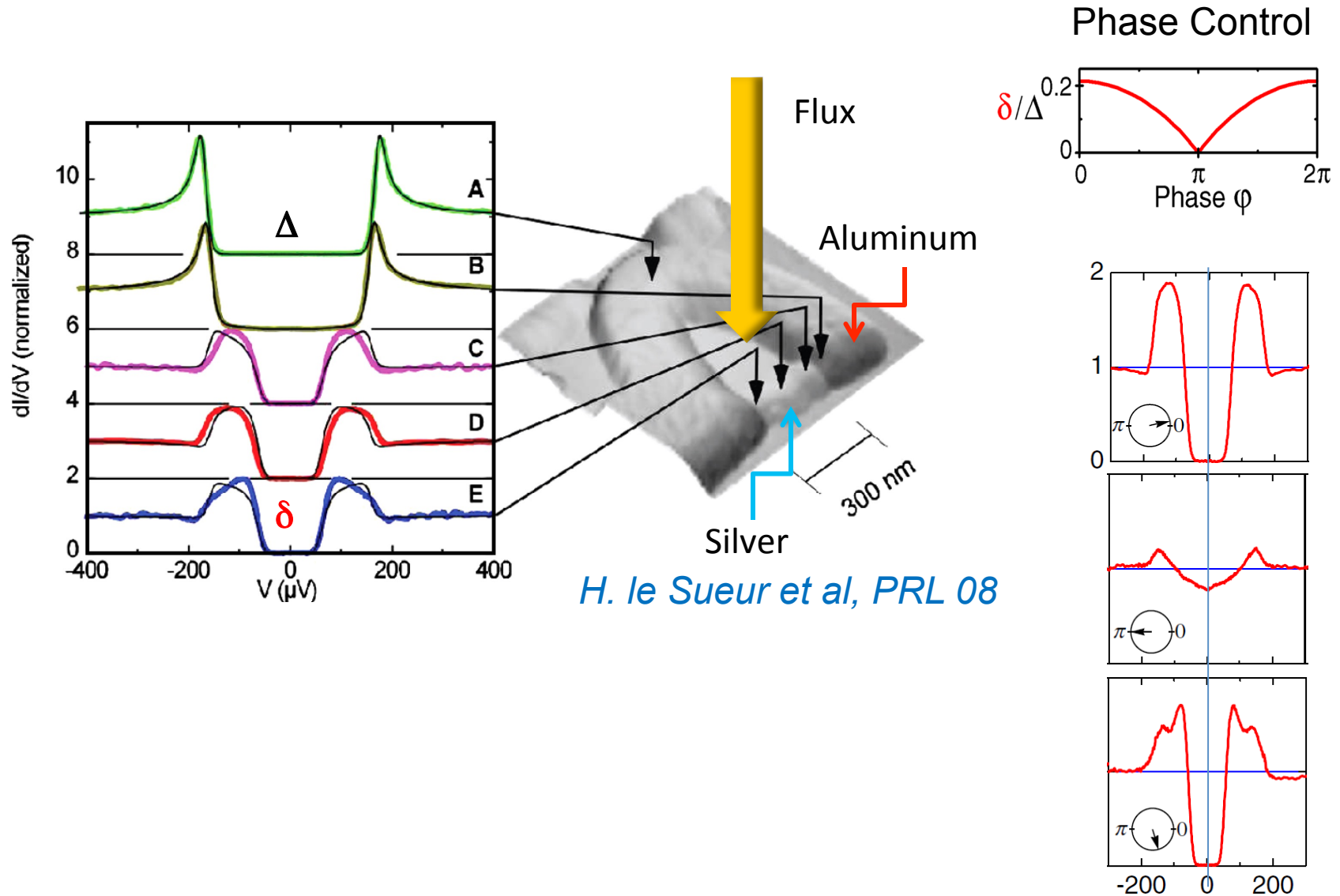
Mi et al, PRB 13

Magnetic Impurity and Supercurrent Isolation



- Gap out topological wire by supercurrent gradient

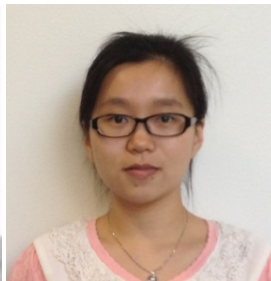
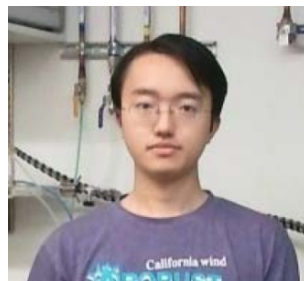
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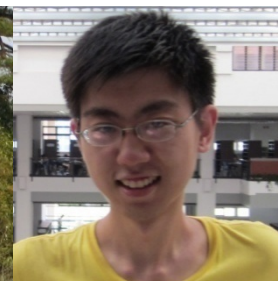
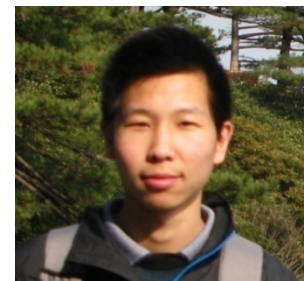
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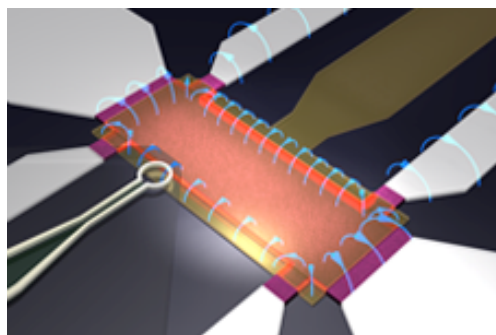
Gerry Sullivan



MBE

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Kathryn Moler
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SQUID Imaging

Theory:

Carol Beenakker

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Dmitry Pikulin

Shoucheng Zhang

.....

Among many others

SUMMARY

- 1) Robust Quantum Spin Hall Effect:
current technology 77K
- 2) InAs/GaSb Helical Edge: clean and highly-tunable
1D system limited only by intrinsic factors
- 3) InAs/GaSb Edge + Superconductor:
a realistic route towards Majorana bound State

THANKS